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1966-53

L. P. Whitehill

Survey on Earth Albedo

11 October 1966

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MASSACHUSETTS INSTITUTE OF TECHNOLOGY

Lexington, Massachusetts



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# MASSACHUSETTS INSTITUTE OF TECHNOLOGY LINCOLN LABORATORY

## SURVEY ON EARTH ALBEDO

L. P. WHITEHILL

Group 63

TECHNICAL NOTE 1966-53

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#### ABSTRACT

This paper is a study of the conditions that can be expected for the intensity of light reflected and scattered from the earth. The albedo of the earth is defined and the albedo of various surfaces, and regions of the earth are discussed. An attempt is made to survey typical results of measurements of Tiros data. The albedo of the earth covered by a Rayleigh atmosphere is presented in graphical form and the intensity distribution over the earth is plotted for various orientations of observer viewed spot and sun. A calculation of expected minimum response is made for a specific photodiode for a certain geometry and a method is presented to enable interested persons to easily perform similar calculations for different geometries.

Accepted for the Air Force Franklin C. Hudson Chief, Lincoln Laboratory Office

## SURVEY ON EARTH ALBEDO

Before beginning the discussion, the oft-used term "albedo" must be clearly understood. The albedo of an object is simply the ratio of the emergent flux of light scattered and reflected in a given spectral region to the incident flux of light in the same spectral region. It is mathematically computable if the emergent spectral intensity [in units of (power)(unit projected area)<sup>-1</sup> (steradian)<sup>-1</sup> (unit wavelength interval)<sup>-1</sup>] is known. Then

$$\mathbf{a}_{1} = \mathbf{a}_{1} (\lambda_{1}, \lambda_{2}) = \text{albedo} = \frac{\int_{\lambda_{1}}^{\lambda_{2}} \int_{0}^{2\pi} \int_{0}^{\pi/2} \mathbf{I}_{\lambda} (\theta, \varphi) \cos \theta \sin \theta \, d \theta \, d \varphi}{\int_{\lambda_{1}}^{\lambda_{2}} \pi \, \mathbf{F}_{\lambda} \cos Z \, d \lambda}$$

where

 $I_{\lambda}$  is the emergent spectral intensity  $F_{\lambda} \text{ is the incident spectral intensity = incident}$  spectral flux[(power)(unit area)<sup>-1</sup> (unit wavelength interval)<sup>-1</sup>] divided by  $\pi$  for a source like the sun.

Z is the solar zenith angle shown in Fig. 1.

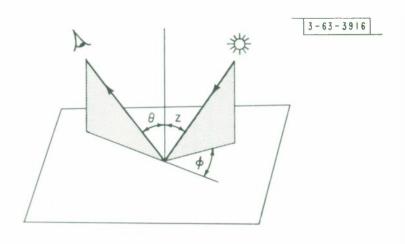


Fig. 1

The albedo just defined is, strictly speaking, the albedo of a point. The albedo of the entire object would have to be obtained by averaging "a" over the object.

To measure the albedo of a point on the earth's surface from an orbiting satellite requires quite a different definition:

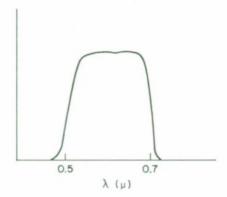
$$a_2 = \frac{F}{\cos Z \int S_{\lambda} \Phi(\lambda) d\lambda}$$

where

F is the flux measured by the sensor  $S_{\lambda} \ \, \text{is the incident spectral flux}$   $\Phi(\lambda)$  is the effective spectral response of the sensor

The value F is not directly measureable because, to accurately measure F, sensors would have to be placed at all points in a hemisphere centered on the viewed area and the intensities would have to be averaged over the hemisphere. Since this is clearly not a practical idea, an assumption has to be made as to the scattering function of the area. The assumption that is invariably made is that the surface scatters isotropically, i.e., according to Lambert's law. It is then a simple matter to produce a number  $F = \pi I$  since the intensity in all directions is assumed constant.

The albedo of an area measured in this way is dependent (especially for a portion of the earth where scattering from the atmosphere is the major contributor) on the exact shape of the spectral response curve of the sensor. To see this assume two sensors with their 3 db points at the same wavelength but having differently shaped spectral response curves as in Fig. 2 below. If the relative intensity of the scattered sunlight as function of wavelength is as shown below (Fig. 3), sensor (b) will record a much higher albedo than sensor (a).



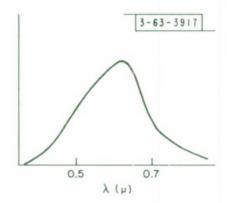


Fig. 2

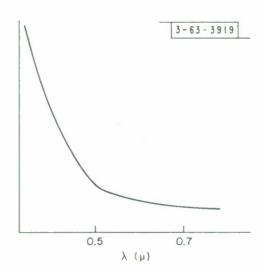


Fig. 3

The realization that the measured albedo depends on the exact shape of the spectral response curve as well as the bandwidth has never been encountered in the literature.

Two alternative definitions of albedo are proposed below, although they are not definitions of albedo in the generally accepted sense and would more properly be called (for want of a better name) backscatter coefficients.

$$b_1(\lambda) = \frac{I_{\lambda}(\theta, \varphi, Z)}{H_{\lambda}} = \text{spectral backscatter coefficient}$$

where

 $\mathbf{I}_{\lambda}$  is the spectral intensity of the scattered sunlight at a given wavelength

 $H_{\lambda}$  is the spectral intensity of sunlight at the same wavelength

Since only the ratio of these two factors is important,  $H_{\lambda}$  may be set to l for theoretical calculations. The spectral flux seen through a solid ''dw'' from a point above the surface is just  $b_1 H_{\lambda} dw$ .

$$b_2 = \frac{\int_{\lambda} \frac{I_{\lambda}(\theta, \varphi, Z) \Phi(\lambda) d\lambda}{\int_{\lambda} H_{\lambda} \Phi(\lambda) d\lambda} = \text{sensor backscatter coefficient}$$

where  $\Phi$  is the effective spectral response of the sensor.

The effective flux flowing into "dw" is just

$$b_2 \int_{\lambda} H_{\lambda} \Phi(\lambda) d\omega = dS(\theta, \varphi, Z)$$

These backscatter coefficients are explicit functions of 3 variables  $\theta$ ,  $\phi$ , and Z, and implicit functions of the type of area viewed, i.e., forests, deserts, oceans, etc., and the sensor itself. Hence, they would be very difficult to tabulate. All that is undertaken in the latter half of this paper is to obtain values for  $b_1$  and  $b_2$  for very specific minimum intensity cases for a specific sensor.

<sup>\*</sup> The Fairchild Planar Photodiode 1N3734

The following is a survey of some of the information available on albedo of the earth as it was defined in this paper.

The albedo of the entire earth is generally computed by observing earth-light reflected from a new moon or by satellite measurements. In the visible region 30% - 40% are respectable limits for the albedo of the earth at any time with a time average value of 36-39% depending on the investigator. For a cloudless earth the albedo is 16%, for a cloud-covered earth it is 56%. The albedo is about 50% in the UV region and 28% in the IR region.

The albedo of various cloud types for the entire spectrum is shown in Table I below.  $^{\rm l}$ 

## TABLE I

Albedo of Various Cloud Types.

Cloud Type	Albedo
Very dense clouds of extensive area and great depth	78
Dense clouds, quite opaque	55-62
Dense clouds, nearly opaque	44
Thin clouds	36 - 40
Stratocumulus, overcast	56 - 81
Altostratus, occasional breaks	17-36
Altostratus, overcast	39-59
Cirrostratus and altostratus	49-64
Cirrostratus, overcast	44-50
Stratus, 680-1000 ft. thick	78

Data for several types of terrain are given below in Table II.

TABLE II

Albedo of Various Surfaces.

Surface	Albedo (per cent)
Desert	24-28
Fields, various types	3-25
Forest, green	3-10
Grass, various conditions	14-37
Ground, bare	7-20
Sand, dry	18
Sand, wet	9
Snow or ice	46 - 86
Water (direct sun only) Z (degrees)	
0	2.0
20	2.1
40	2.5
50	3.4
60	6.0
70	13.4
80	34.8
85	504
	58.4

where Z is the zenith angle of the sun.

The albedo of water according to one author is closely approximated by the above numbers, calculated from Fresnel's law of reflection, even when the wind is 10 mph. Angstrom states that in ordinary geophysical problems the data in Table II are applicable. 1

A good source of values for earth albedo is the results of radiation measurements from the Tiros meteorological satellite series. The results gotten by Nordberg and Bandeen are given below as an example of what information can be culled from Tiros data.

The Tiros satellites are generally equipped with two types of radiometers. The first is a non-scanning radiometer with a broad response in both the visible and infrared regions and a rather low spatial resolution. Its aperture is approximately 55 degrees. The second is a scanning radiometer which scans as the satellite spins. It is of medium spatial resolution, has a 5 degree aperture, and responds to radiation in five different spectral regions separated by filters. Three of these regions are in the infared between 5.0 and 6.7 microns, between 7.5 and 13.5 microns, and between 7.0 and 32.0 microns. The other two regions lie mainly in the visible portion of the spectrum, that is, between 0.2 and 7.0 microns and between 0.50 and 0.75 microns. It is the latter, channel 5, which is of interest here. Note: The satellite is flying in a nearly circular orbit with an approximate altitude of 780 km.

Shown in Fig. 4 through Fig. 6 are albedo maps made with the channel 5 sensor. These values of albedo, as all other albedo values derived from Tiros data, assumes an isotropic reflecting surface.

The data of interest for minimum intensity considerations is to the left of the sub-satellite path where specular reflection from the water does not occur. Minimum albedo for this area, with the sun near zenith, is 4%.

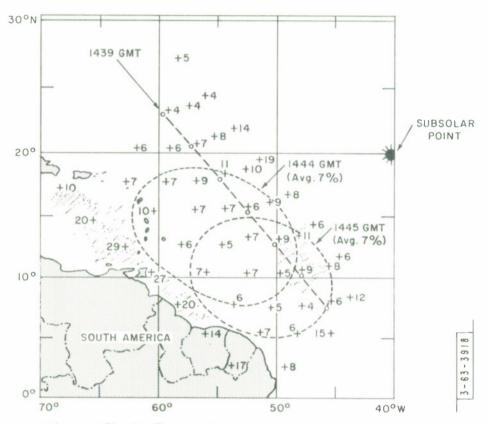


Figure 4. Tiros Albedo Percentages

Albedo percentages measured by channel 5 (approximately  $0.5 - 0.75\mu$ ) of the medium resolution radiometer between 1439 and 1445 GMT, Orbit 117, 20 July 1961. The skies are mostly clear, except for some minor scattered clouds. Averages of these data are given for the areas enclosed by dashed lines. The heavy dashed line is the sub-satellite path with points each minute.

The next albedo map is for the eastern United States under mostly cloudy skies.

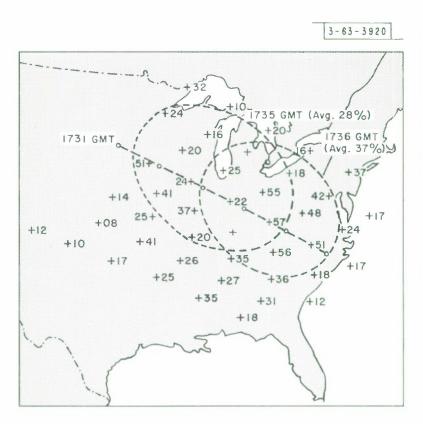


Figure 5. Albedo Percentages Measured by Channel 5.

Notice the low value of 84 in the absence of clouds.

The last albedo map shows the North African Desert under clear skies.

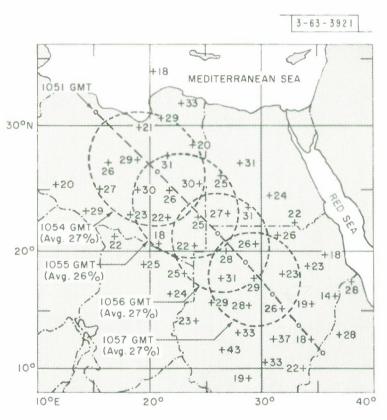


Figure 6. Albedo Percentages Measured by Channel 5.

The albedo values presented in the above maps do not take into account the fact that the earth does not scatter isotropically. To get a quantitative idea as to the extent to which anisotropy occurs, the following computation with Tiros III data was performed. <sup>3</sup>

A selected area of the earth, Fig. 7, is viewed by the radiometer twice during the orbit with the scans 3 to 4 minutes apart. The satellites have moved far enough in the time for the scattering angle of the scattered and reflected radiation to have changed quite significantly. The two maps above show the same area. Notice the difference in the appearance of the two maps, a difference which is not acknowledged in the previous albedo maps. It should be noted that the smaller scattering angles in this study are associated with the larger nadir angles. There is a large difference in nadir angle between the two sets of data along X. However, when this difference is reduced, as along Y, the fluxes for the smaller scattering angles remain considerably higher. Since the solar zenith angle varies in an identical manner along both lines, the difference in flux can be least qualitatively be attributed to the difference in scattering angle. The variation of channel 5 flux along X and Y for the two different viewing geometry is shown in Fig. 8.

The data of both solar channels suggest that the scattering angle effect is largest in areas of least cloudiness. This suggests the possibility that specular reflection from the ocean surface could have been responsible for the larger flux differences in the areas of relatively low flux. Along Y, this effect could have entered since this line lies nearly in the plane of the sun. Since both X and Y show the same relationship between flux difference maximas and regions of relatively low flux, specular reflection alone could not have been the cause. Increased atmospheric scattering for the smaller scattering angles, regardless of the viewed surface, is probably the important factor.

In the report just considered, the data sample included significant differences in the radiometric nadir angle associated with the separate scans of the overlapping data. Therefore, the increase in forward scatter could have been effected by the variation in the viewed volume. However, even as in a subsequent publication, when the nadir angle difference of the two scans is reduced to zero, the measured albedo remain higher for smaller scattering angles (Fig. 9).

Notice in the figure that in general, effective blackbody temperature measured with the channel 2 sensor (approximately 8.0-12.0 microns) decreases

<sup>\*\*</sup>scattering angle = 180° -(angle between sun, viewed spot and satellite sensor)
\*\*nadir angle = angle between local vertical at the sensor and the direction of the viewed spot.

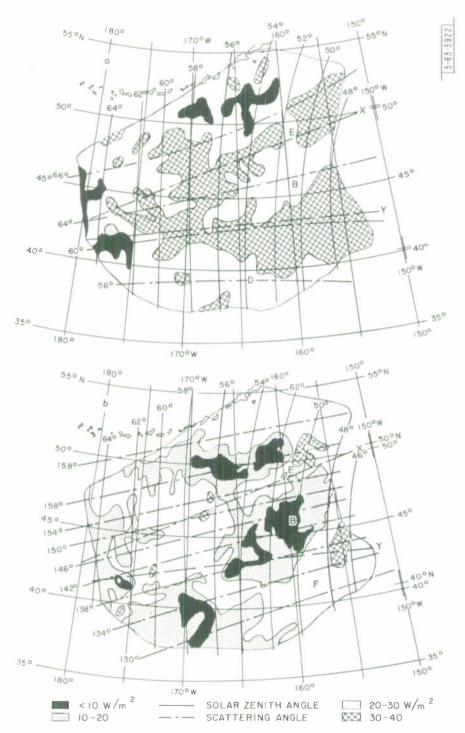
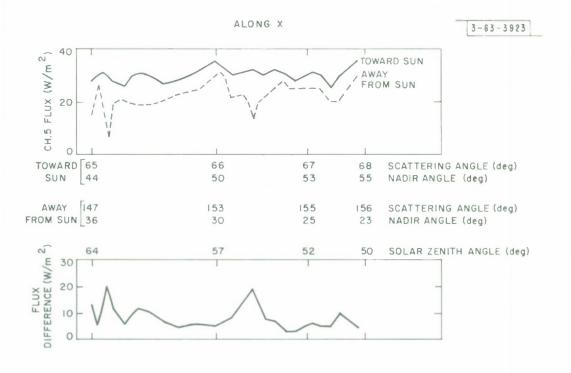


Figure 7. Relative Patterns of Channel 5 Flux for Area with Overlapping Scanlines from Orbit 5, Tiros III

- (a) Initial Alternating Open Mode Data
- (b) Complemental Single Open Mode Data



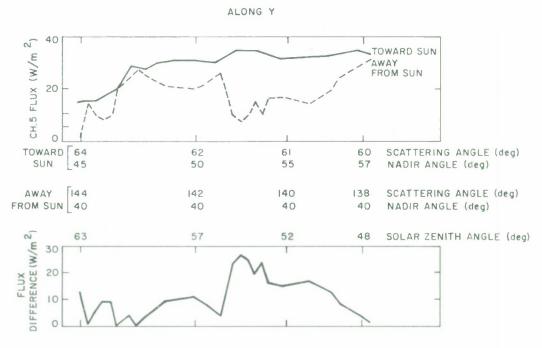


Figure 8. Variation of Channel 5 Flux for Radiometer Scans Towards and Away from the Sun and the Variation of Their Difference Along X and Y.

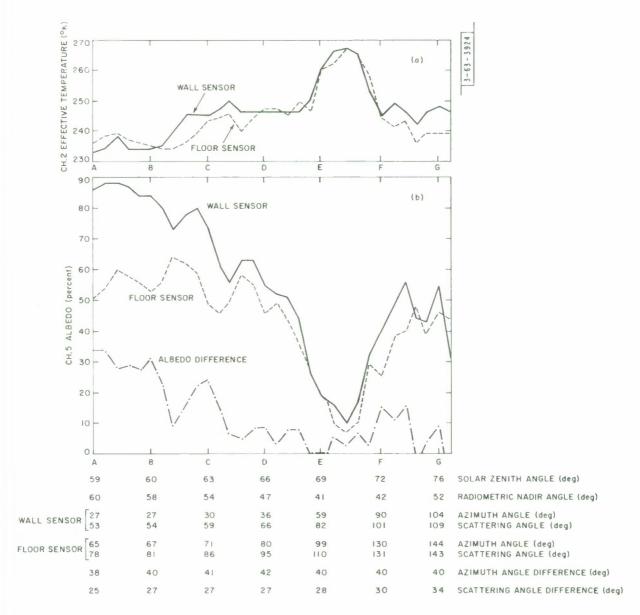


Figure 9. Comparison of Data from the Wall and Floor Sensors Along the Isopleth of Zero Nadir Angle Difference

(a) Channel 2 Effective Temperatures

(b) Channel 5 Albedos and Albedo Differences

as the albedo increases. This is due to the fact that the tops of clouds are much colder than the less reflective underlying surface.

It would seem that a careful study of raw Tiros data could yield information on the minimum light intensities and reflected and scattered from the earth. This path of investigation was not chosen, however, because of two considerations. The spectral regions of the Tiros sensors and the sensor under consideration are different enough to make the Tiros data almost inapplicable for this study; also it has been noticed that many authors tend to disregard Tirosderived albedos whenever they fall below the albedo that can be expected from scattering by a Rayleigh atmosphere. As an example one author found the albedo of a point over ocean to be 6%. For the same geometry and for approximately the same spectral regions the albedo of a standard atmosphere with Rayleigh scattering and zero surface reflectivity deduced by computations by K. L. Coulson is 8%. It is stated that the difference is dependent on the degradation of the instrumental response, although a proper evaluation should include the effects on the Rayleigh albedo of molecular absorption and the scattering and absorption by aerosols.

#### RAYLEIGH SCATTERING

The problem of atmospheric scattering is not an easy one and has been tackled only recently by Chandrasekhar in his book on Radiative Transfer. <sup>5</sup>

The equations which he introduces have been tabulated by K. L. Coulson in his various publications on scattering by a Rayleigh atmosphere. Some of his results and the results of others working along similar lines are discussed in the succeeding pages.

The values that follow have been computed for parallel unpolarized radiation of net flux  $\pi$  per unit area incident on the top of an idealized plane parallel atmosphere. The atmosphere is assumed to be uniform and infinite in the horizontal, non-absorbing, composed of scattering centers which are small compared with the wavelength of the light, and of finite optical thichness.

For the sun at a given zenith angle  $\theta_0$  the total outward flux from the top of the atmosphere in a unit frequency interval and unit time is given by

$$D_{v} = \int_{0}^{2\pi} \int_{0}^{\pi/2} I_{v}(\theta, \varphi) \sin \theta \cos \theta d \theta d \varphi$$

 $\theta$  ,  $\phi$  , and  $\theta$   $_{0}$  are defined by Fig. 1 with Z replaced by  $\theta$   $_{0}.$   $I_{_{V}}$  ( $\theta$  ,  $\phi$ ) is a tabulated function.

The spectral albedo or the monochromatic reflectance is

$$R_{v} = \frac{D_{v}}{\pi \mu_{o}}$$

where

$$\mu_0 = \cos \theta_0$$

Below,  $R_{y}$  is tabulated for various wavelengths and zenith angles.

Table III is represented graphically in Fig. 10. The values indicate the reflectance increases for decreasing wavelength at all sun elevations due to the inverse fourth power dependence of Rayleigh scattering. <sup>5</sup>

TABLE III

Reflectance of the Atmospheric Model at Various Wavelengths for Representative Values of Sun Elevation and Surface Reflectivity. (Lambertian reflector)\*6

A	μ 0	λ3150A	4000	5000	6000	7000	8000
0	0.02	0.754	0.638	0.565	0.503	0.430	0.325
0	0.10	0.705	0.552	0.396	0.247	0.154	0.094
0	0.40	0.550	0.304	0.152	0.079	0.044	0.024
0	0.80	0.400	0.181	0.085	0.039	0.022	0.013
0	1.00	0.345	0.148	0.068	0.033	0.017	0.009
0. 25	0.02	0.791	0.712	0.665	0.625	0.574	0.494
0.25	0.10	0.752	0.643	0.536	0.427	0.362	0.319
0.25	0.40	0.623	0.445	0.345	0.298	0.276	0.285
0.25	0.80	0.492	0.347	0.293	0.271	0.260	0.257
0.25	1.00	0.450	0.322	0.280	0.264	0.256	0.255
0.80	0.02	0.922	0.910	0.904	0.896	0.885	0.865
0.80	0.10	0.918	0.892	0.866	0.842	0.826	0.815
0.80	0.40	0.860	0.830	0.812	0.805	0.803	0.800
0.80	0.80	0.811	0.798	0.798	0.798	0.798	0.799
0.80	1.00	0.794	0.792	0.794	0.796	0.798	0.799

<sup>\*</sup> Note a more complete table by Coulson is given in Appendix I.

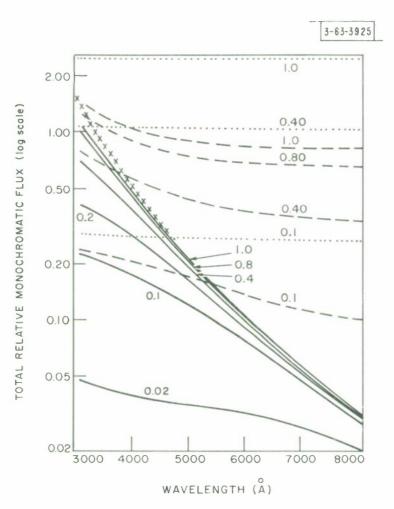


Figure 10. Total relative monochromatic flux emerging from the top of the atmosphere as a function of wavelength, for various sun elevations and three values of surface reflectivity: A = 0 (solid curves), A = 0.25 (dashed curves) and A = 0.80 (dotted curves). Curves are labeled with values of  $\mu$ . A curve (X) for primary Rayleigh, scattering and no attenuation is included for comparison.

This graph, however, is only relative and assumes all the incident flux to be the same at all wavelengths. If we assume a source with the same variation with wavelength as the sun the graph of Fig. ll is obtained.  $^6$ 

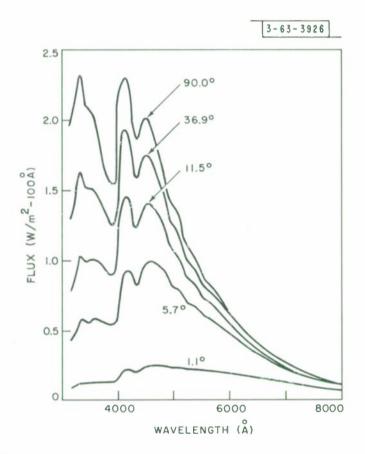


Figure 11. Absolute monochromatic flux from the top of the model atmosphere as a function of wavelength, for zero surface reflectivity and five different sun elevations ( $\mu_0$  = 1.0, 0.40, 0.20, 0.10, 0.02; A = 0)

Using the graphs of Fig. 11 the albedo for a plane parallel earth as a function of sun elevation for the spectral region 3200 A  $\leq \lambda \leq$  8000 A can be obtained. The results obtained by Coulson are plotted in Fig. 12 and tabulated in Table IV.

TABLE IV

Values and Comparisons of the Total Absolute Flux for 3200 A  $\leq \lambda \leq 8000$  A at Six Different Sun Elevations. (Units of flux: W/m<sup>2</sup>)

Sun elevation	1.1°	5.7°	11.5°	23.6°	53.1°	90.0°
Fo	15.4	77.0	154.0	307.9	615.8	769.9
(J)	8.0	25.0	34.3	44.0	48.9	51.8
(J) <sub>25</sub>	9.7	37.1	62.4	103.6	179.9	214.8
(J) <sub>80</sub>	13.7	65.9	128.3	249.6	488.6	609.6
$(\alpha)_0$	0.519	0.325	0.223	0.143	0.079	0.067
$(\alpha)_{2.5}$	0.630	0.482	0.405	0.336	0.292	0.279
(α) <sub>80</sub>	0.890	0.856	0.833	0.811	0.793	0.792
(F <sub>0g</sub> ) <sub>0</sub>	7.4	52.0	119.7	263.9	566.9	718.1
(F <sub>0g</sub> ) <sub>25</sub>	7.6	53.2	122.1	272.4	581. 2	740.1
(F <sub>0g</sub> ) <sub>80</sub>	8.5	55.5	128.5	291.5	636.0	801. 5

\* F : Incident flux on horizontal surface at top of atmosphere.

(J)<sub>A</sub>: Emergent flux through horizontal surface at top of atmosphere. (Subscript A represents value of surface reflectivity: 0.025 or 0.80).

 $(\alpha)_A$ :  $\frac{(J)_A}{F_0}$  = Albedo: ratio of emergent to incident flux.

 $(F_{0\sigma})_A$ : Flux incident on horizontal surface at ground.

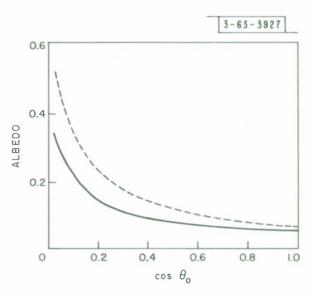


Figure 12. Approximated albedo for the entire solar spectrum (solid curve) and computed albedo for the spectral range  $3200 \text{ A} \ge \lambda \ge 8000 \text{ A}$  (dashed curve) as a function of sun elevation (A = 0).

Coulson also did a calculation of the planetary albedo of the spherical earth due to atmospheric scattering. The disk of the earth was divided into concentric plane parallel rings and the sun elevation for each ring was considered constant. A value of 7.6% was obtained for the entire spectrum. This value is reduced by ozone absorption in the region near and below 3000 Å. The corrected albedo for this more realistic case is given as 6.9%.

Some additional graphs computed by Coulson in a later publication for a plane parallel atmosphere with a Lambertian surface at the bottom are presented in Figs. 13-17. Figures 13-15 are a graphical representation of the values tabulated in Appendix I. It can be easily seen that the atmosphere itself scatters back a large portion of the short-wavelength radiation especially for the larger solar zenith angles, and that at longer wavelengths and smaller solar zenith angles the effect of the reflectance of the underlying surface becomes predominant. Figures 16-17 show the strong dependence of system reflectance for large solar zenith angles particularly pronounced for small values of A (for the entire solar spectrum (Fig. 16) and for 3 spectral regions (Fig. 17)).

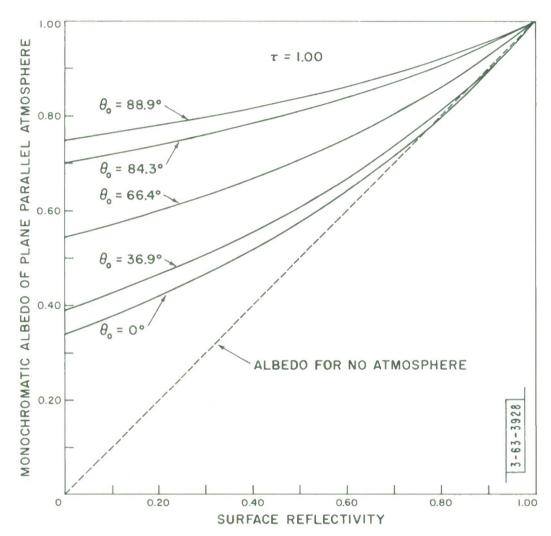


Fig. 13 Monochromatic reflectance of model atmosphere-plus-surface system as a function of surface reflectivity for optical thickness  $_{\text{T}}$  = 1.00 for five different solar zenith angles  $_{\text{T}}$  = 1.00, corresponding to wavelength  $_{\text{A}}$  = 3120 Å at sea level.

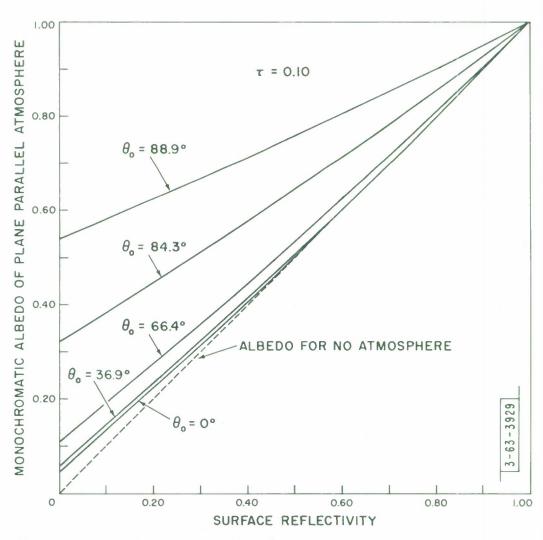


Fig. 14 Monochromatic reflectance of model atmosphere-plus-surface system as a function of surface reflectivity for optical thickness  $\tau$  = 0.10 for five different solar zenith angles  $\tau$  = 0.10, corresponding to wavelength  $\lambda$  = 5460 Å.

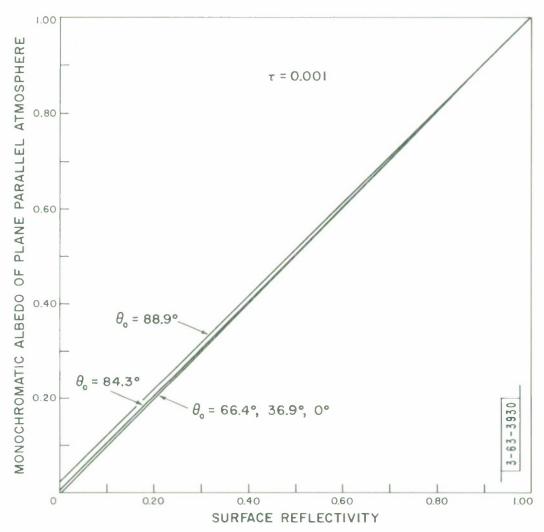


Fig. 15 Monochromatic reflectance of model atmosphere-plus-surface system as a function of surface reflectivity for optical thickness  $_{\text{T}}$  = 0.00l for five different solar zenith angles  $_{\text{T}}$  = 0.00l, corresponding to wavelength  $_{\text{T}}$  = 17,000Å.

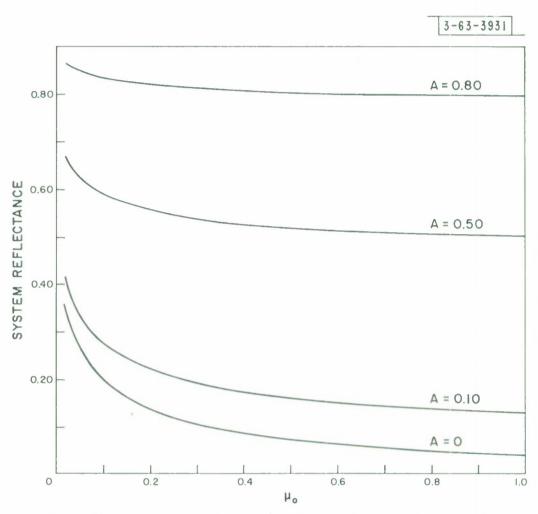


Fig. 16 Reflectance of model atmosphere-plus-surface system as a function of cosine of solar zenith angle  $\mu$  for spectral interval 3100 Å  $\leq \lambda \leq$  25,000 Å for different values of surface reflectivity.

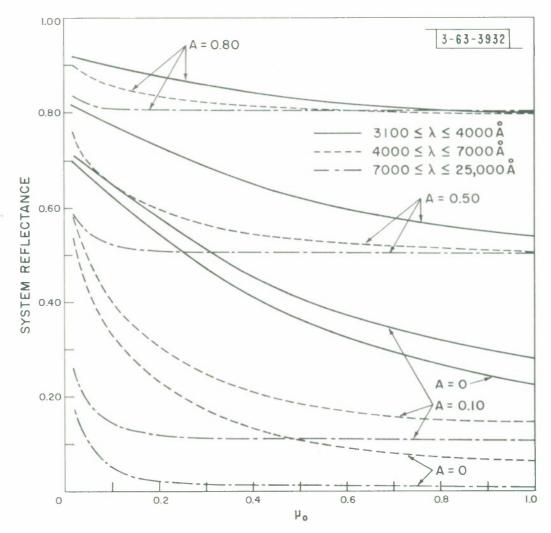


Fig. 17 Reflectance of model atmosphere-plus-surface system as a function of cosing of solar zenith angle  $\mu$  , with spectrum divided into ultraviolet (3100 Å  $<\lambda<4000$  Å), visible (4000 Å), and infrared (7000 Å  $<\lambda<25,000$  Å)

The graphs and tables that have been presented so far, for the most part, have indicated the total amount of light reflected from the plane-parallel atmosphere. Relatively little information has been presented about the expected minimum light conditions at various configurations of observer, viewed spot and source. The material on the succeeding pages will provide information on the expected intensities and "backscatter coefficients" of the reflected and scattered light.

Some of the results are presented in a form which makes it necessary to discuss briefly the concept of optical thickness,  $\tau$ .  $\tau$  is defined as

$$\tau = \int_{0}^{z_{0}} n(z) \sigma_{\lambda}(z) dz$$

where

n = number of scattering centers/unit volume

 $\sigma_{\lambda}$  = scattering cross-section

 $z_{o}$  is usually taken as the height of the atmosphere

Then the intensity of light on the ground is

$$I = I_o e^{-\tau}$$

 $_{\text{T}}$  is related to  $\lambda$  by the following Table (Table V) and Fig. 18.

TABLE V

1.00	. 50	. 25	. 15	.10	. 05	. 02
3120	2715	4365	4950	5460	6440	8090

All graphs in Figs. 19 and 20 take into account the effect of primary and multiple scattering (unless otherwise indicated) and increased path lengths for oblique rays. The relative intensities indicated on the graphs are relative to an incident intensity of 1.0.

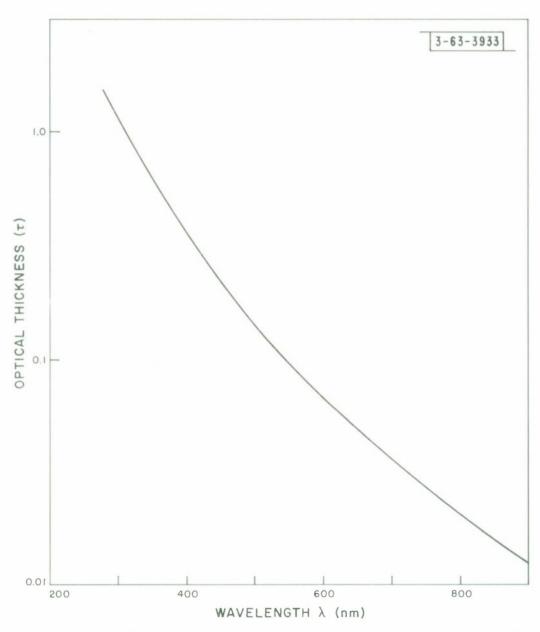


Fig. 18 Optical thickness of the earth's atmosphere as a function of wavelength.

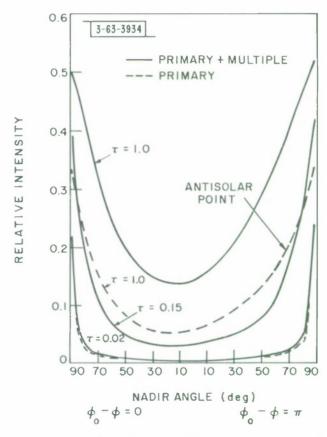


Figure 19. Relative intensity for three different values of optical thickness as a function of direction in the sun's vertical. Dashed curves are primary scattering only; solid curves include higher order scattering ( $\theta_0 = 66.3^{\circ}$ , A = 0).

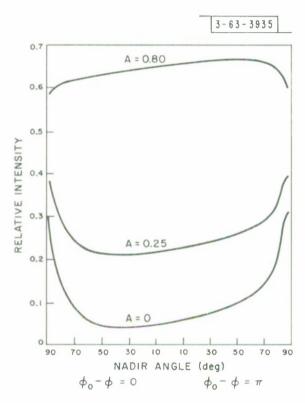


Figure 20. Relative intensity as a function of direction in the sun's vertical, for moderate optical thickness, sun elevation  $53^{\circ}$  and three values of surface reflectivity ( $_{\text{T}} = 0.15$ ,  $\theta_{_{0}} = 36.9^{\circ}$ ).

The plot in Fig. 20 shows that for moderate optical thickness the surface reflected light may constitute the major part of the emergent light. For values of A approaching 1 the surface begins to look like a Lambertian surface. For higher optical thickness the reflectivity is relatively less dominant but very important, even for optical thickness near 1.0.

In a recent paper by William Snoddy <sup>8</sup> the case of the plane-parallel atmosphere was extended for a spherical earth. The portion of the earth visible at a particular altitude was divided into 400 plane parallel sections and the

intensity computed for each section. The irradiation incident on a differential element at a given altitude was computed and compared to the irradiation expected if the earth were a diffuse surface to obtain an "effective albedo". Note that this definition is different from the other definitions of albedo presented earlier and could have extensive use with an earth sensor of very wide aperture. The geometry of the situation is presented in Figs. 21 - 22.

Towards the limb of the earth, incident light travels through a smaller air mass than if the atmosphere were plane-parallel. The air mass through which an incident beam will travel has been calculated by Bemporad and the air mass is given in Fig. 23 together with the air mass for a plane parallel atmosphere. This graph is used to correct the data obtained for a plane parallel atmosphere for large angles of incidence.

Figures 24 - 56 are plots of intensity over the planetary sphere with isophotes drawn through the points of equal intensity. The sun moves to the right along the horizontal diameter of the circles. The differential element is above the center of the circles.

Figures 24 - 30 illustrate the variation as a function of  $\theta_0$  with H =  $10^6$  km,  $\tau$  = 1.0 ( $\lambda$  = 3120 A) and A = 0. Figures 31 - 37 show the same variation with  $\theta_0$  but have H =  $10^3$  km. Variation of intensity with altitude is illustrated in Figs. 38 - 42. Maps showing the effect of the reflectance of the underlying surface are given for an altitude of a million kilometers (Figs. 43 - 45) and a thousand kilometers (Figs. 46 - 48).

Perhaps the most interesting maps from the point of view of using a light sensor with a broad rather than a line bandwidth are the maps showing variation in intensity with optical thickness and hence wavelength. Examples are shown for  $\tau = 1.0$ , 0.5, 0.25, 0.1,  $\theta_0 = 30^\circ$ ; A = 0 and H =  $10^6$  km (Figs. 49-52) and  $10^3$  km (Figs. 53-56).

These relative values of intensity are what has been defined previously as the spectral backscatter coefficient, "b<sub>1</sub>". The quantity  $b_1H_{\lambda}$  dw was then integrated over the sphere and the relative flux intensities plotted in Figs. 57-58. As the reflectance of the underlying surface is increased, the maximum flux intensity shifts away from the ultraviolet towards the maximum of the sun.

The spectral effective albedo of the earth is shown in Fig. 59 for H =  $10^3$  km and three values of A. It should be realized that this graph is valid only for the given altitude. In Figs. 60 - 61 the effective albedo for the spectral region between 3050 and 8050 Å (graphs 1) and for the entire spectrum (graph 2) were calculated as a function of A and H for two values of  $\theta_0$ .

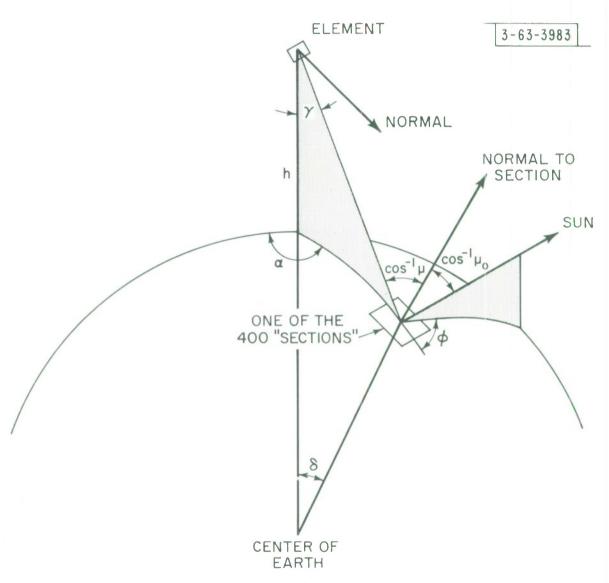


Fig. 21 Section on planet surface

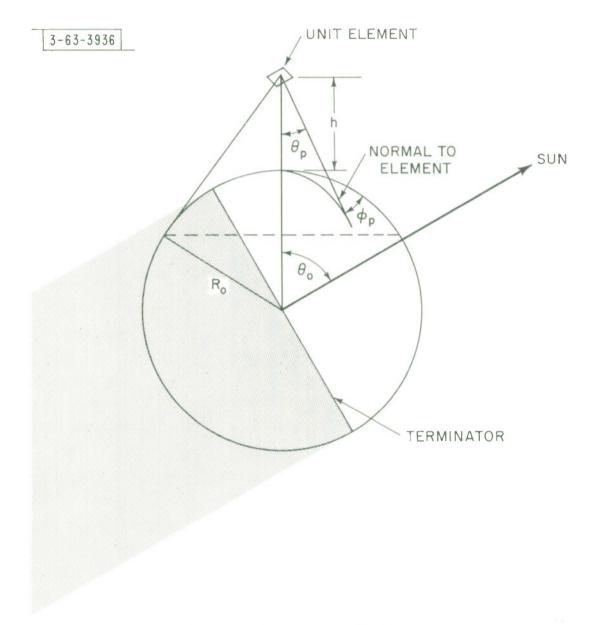


Fig. 22 Element above planet

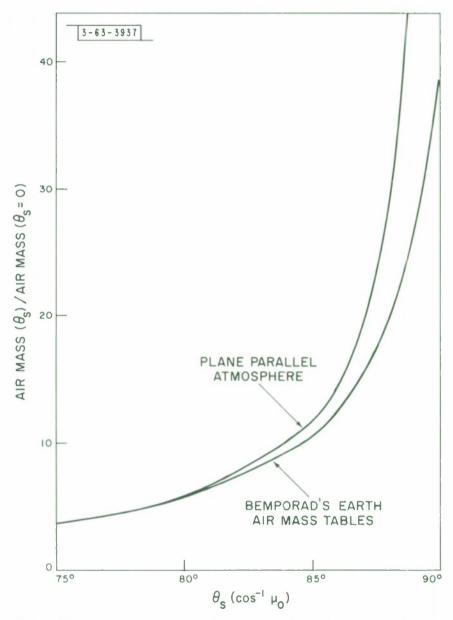


Fig. 23 Air mass as a function of angle of incidence

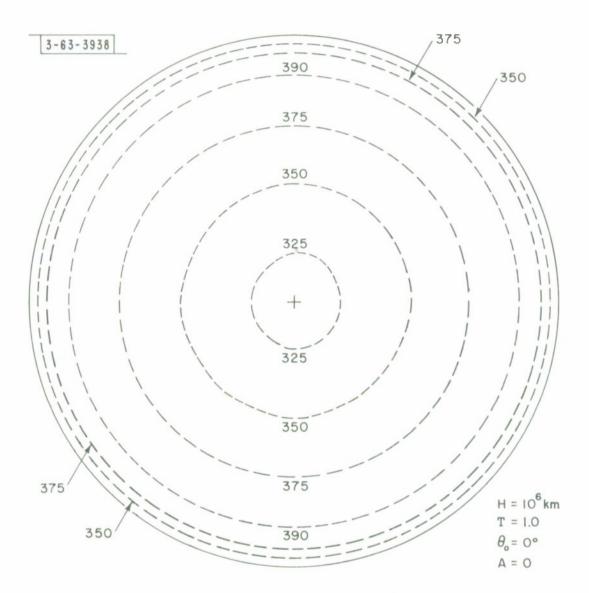


Fig. 24 Map of isophotes for  $H = 10^6$  km,  $\theta_0 = 0$ 

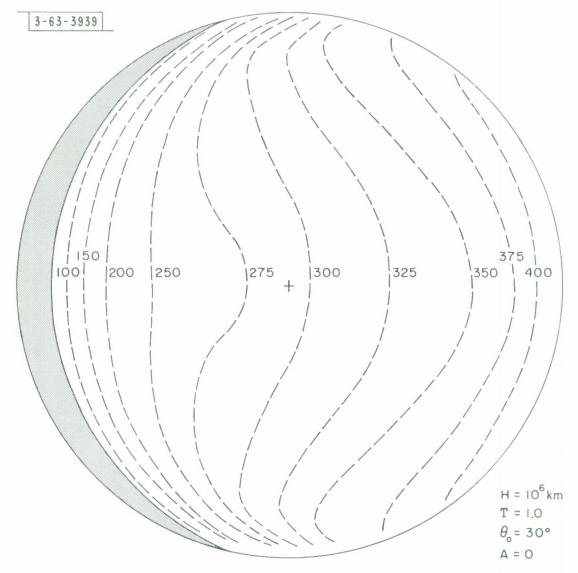


Fig. 25 Map of isophotes for  $H = 10^6$  km,  $\theta_0 = 30^\circ$ 

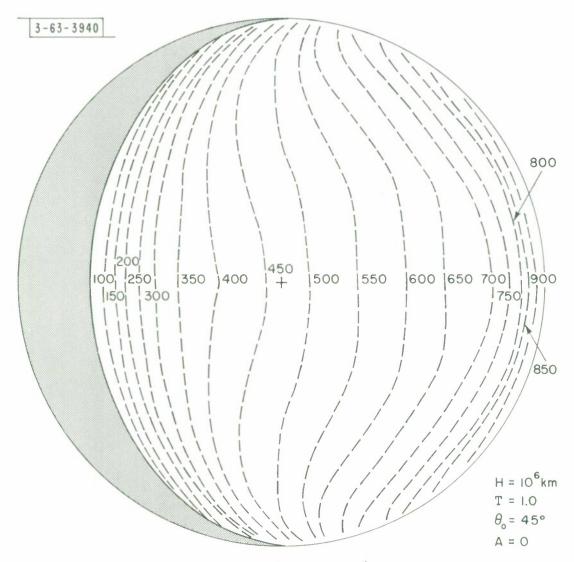


Fig. 26 Map of isophotes for H =  $10^6$  km,  $\theta_0 = 45^0$ 

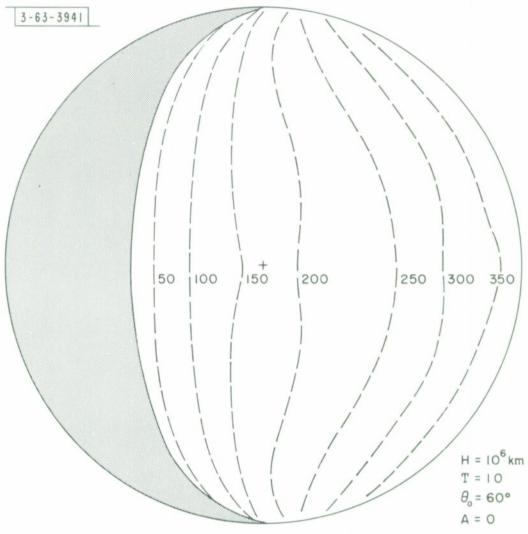


Fig. 27 Map of isophotes for  $H = 10^6$  km,  $\theta_0 = 60^\circ$ 

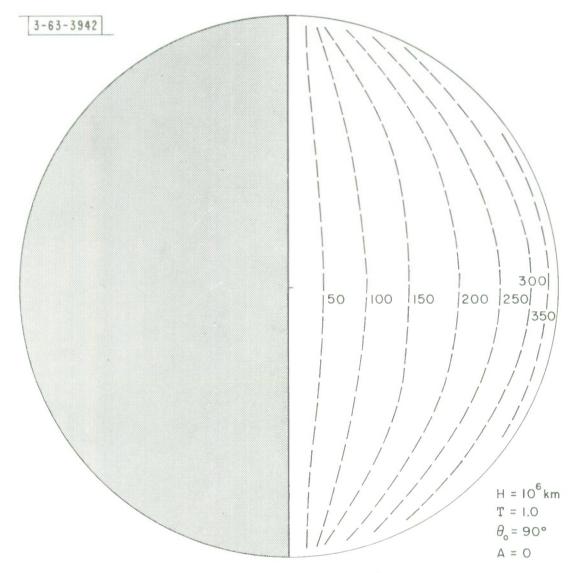


Fig. 28 Map of isophotes for  $H = 10^6$  km,  $\theta_0 = 90^\circ$ 

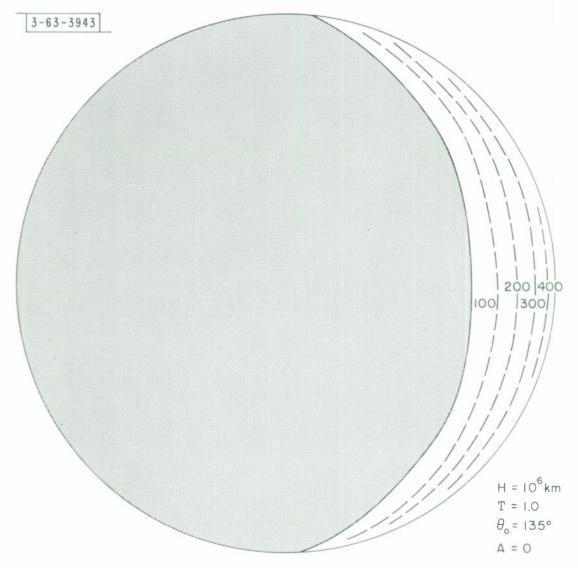


Fig. 29 Map of isophotes for  $H = 10^6$  km,  $\theta_0 = 135^\circ$ 

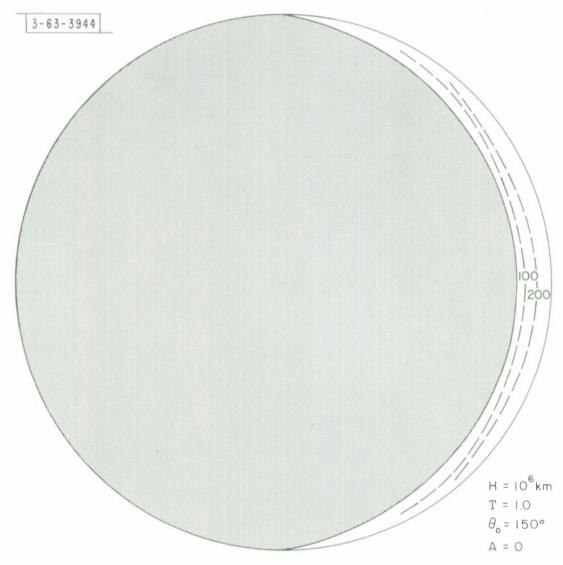


Fig. 30 Map of isophotes for  $H = 10^6$  km,  $\theta_o = 150^\circ$ 

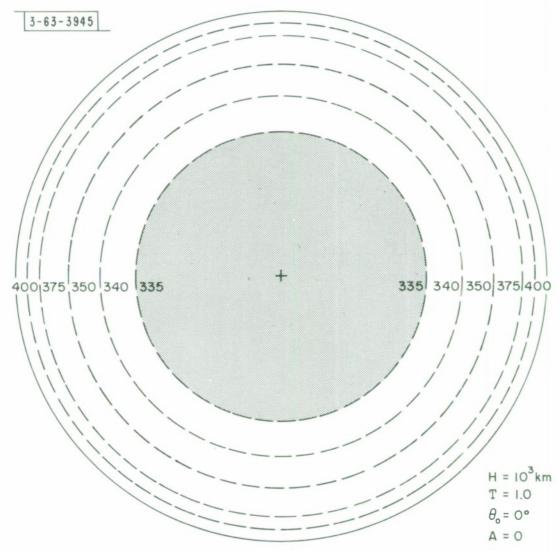


Fig. 31 Map of isophotes for H =  $10^3$  km,  $\theta_0 = 0^0$ 

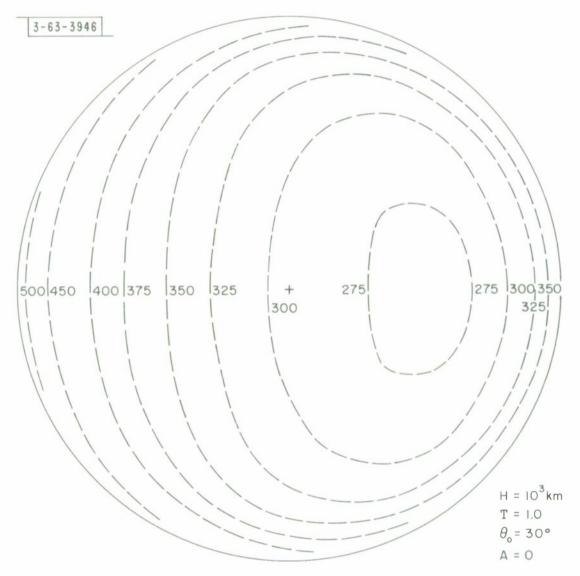


Fig. 32 Map of iosphotes for  $H = 10^3$  km,  $\theta_0 = 30^0$ 

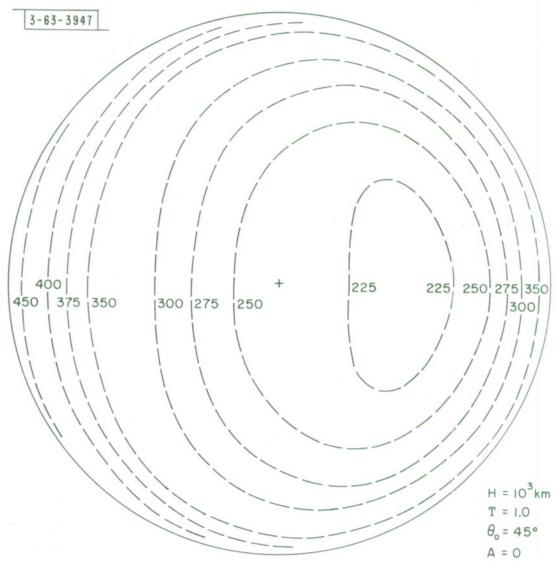


Fig. 33 Map of isophotes for  $H = 10^3$  km,  $\theta_0 = 45^\circ$ 

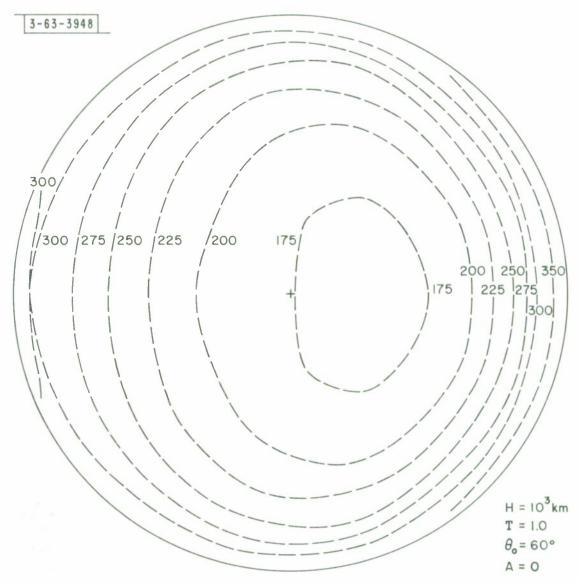


Fig. 34 Map of isophotes for  $H = 10^3$  km,  $\theta_0 = 60^\circ$ 

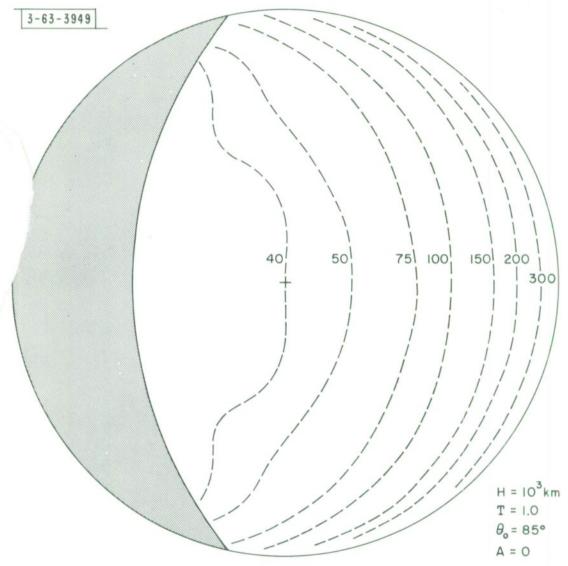


Fig. 35 Map of isophotes for  $H = 10^3$  km,  $\theta_0 = 85^\circ$ 

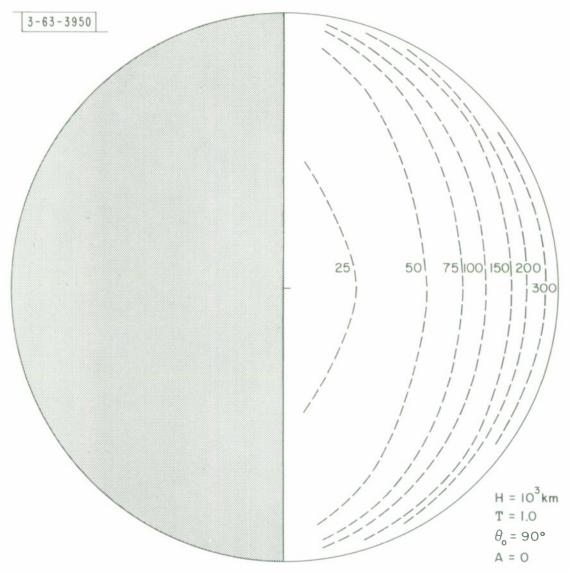


Fig. 36 Map of isophotes for  $H = 10^3$  km,  $\theta_0 = 90^\circ$ 

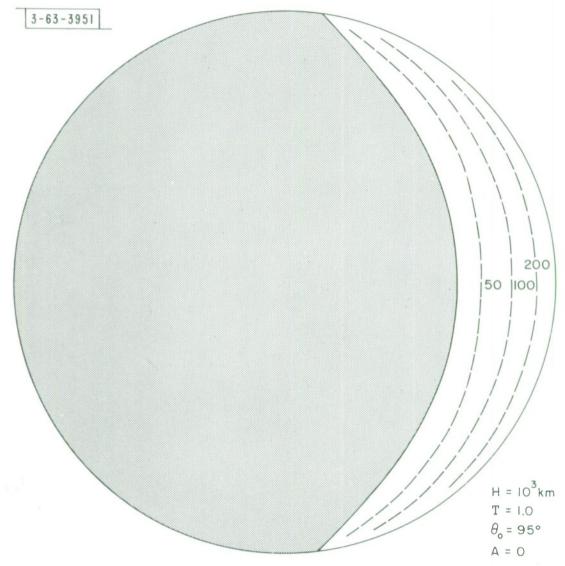


Fig. 37 Map of Isophotes for  $H = 10^3$  km,  $\theta_0 = 95^\circ$ 

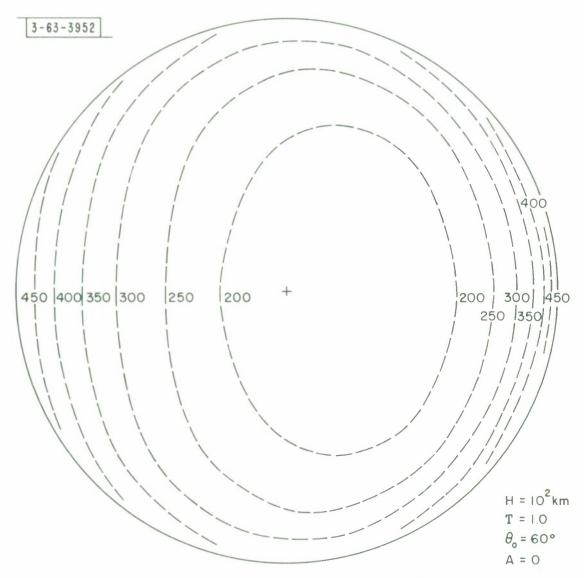


Fig. 38 Map of isophotes for  $H = 10^2 \text{ km}$ 

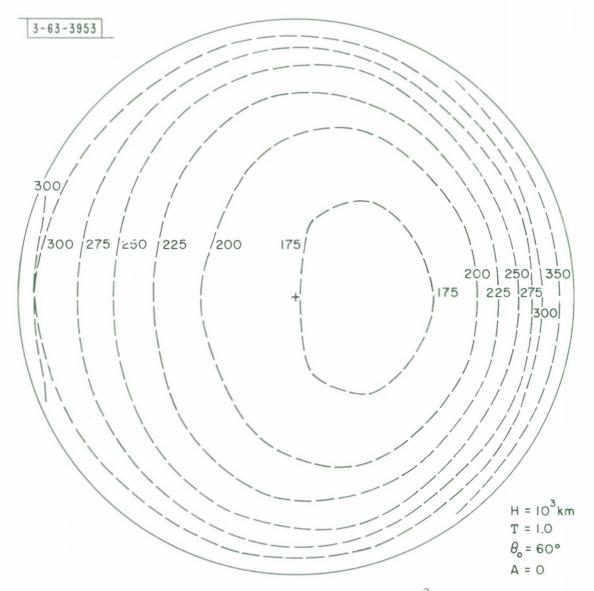


Fig. 39 Map of isophotes for  $H = 10^3 \text{ km}$ 

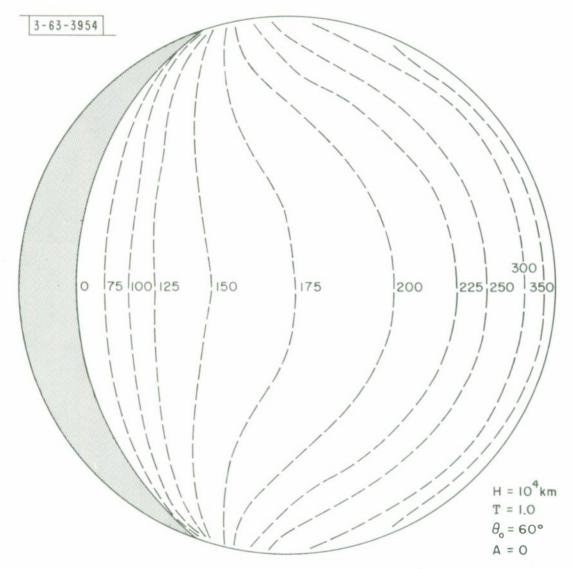


Fig. 40 Map of isophotes for  $H = 10^4$  km

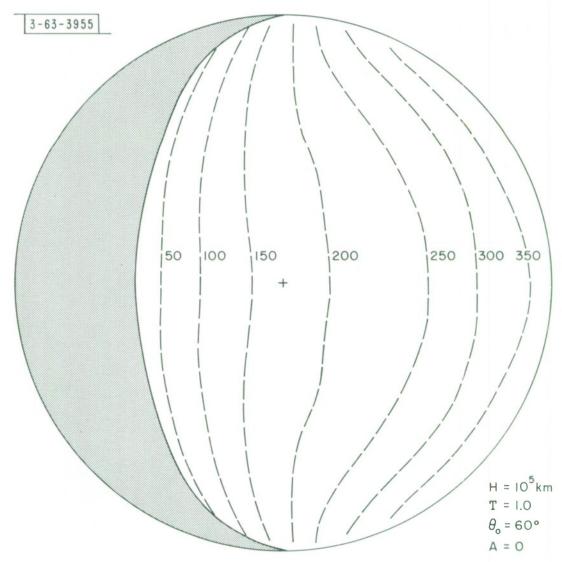


Fig. 41 Map of isophotes for  $H = 10^5 \text{ km}$ 

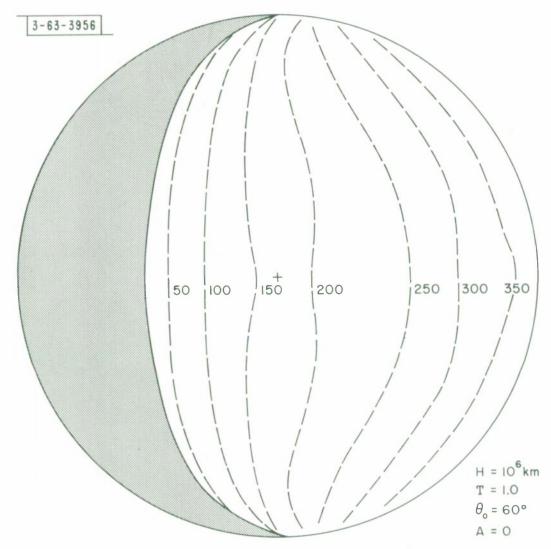


Fig. 42 Map of isophotes for  $H = 10^6 \text{ km}$ 

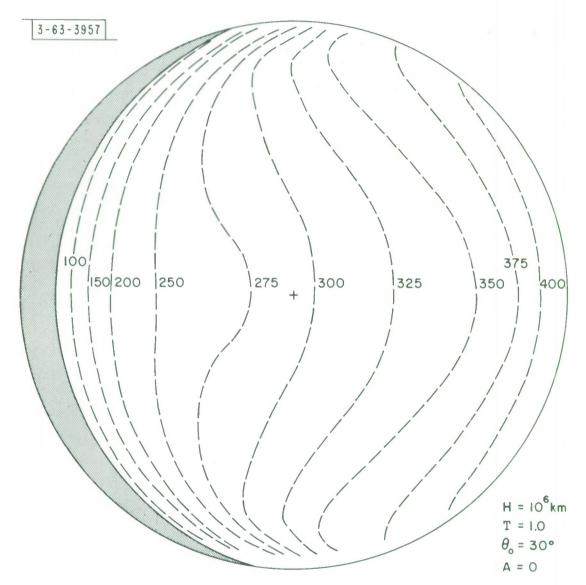


Fig. 43 Map of isophotes for  $H = 10^6$  km, A = 0

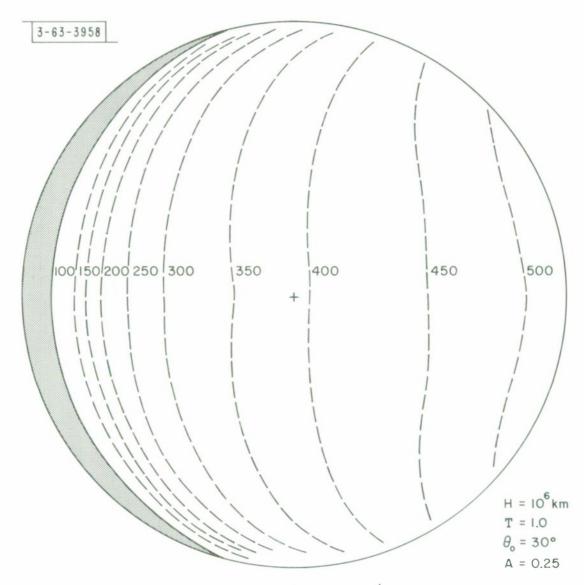


Fig. 44 Map of isophotes for  $H = 10^6$  km, A = 0.25

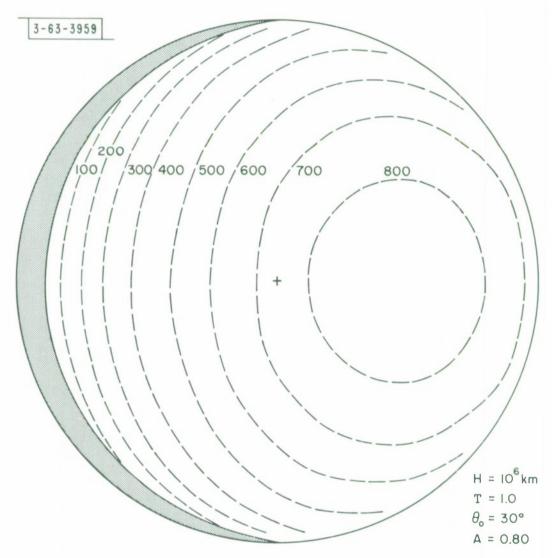


Fig. 45 Map of isophotes for  $H = 10^6$  km, A = 0.80

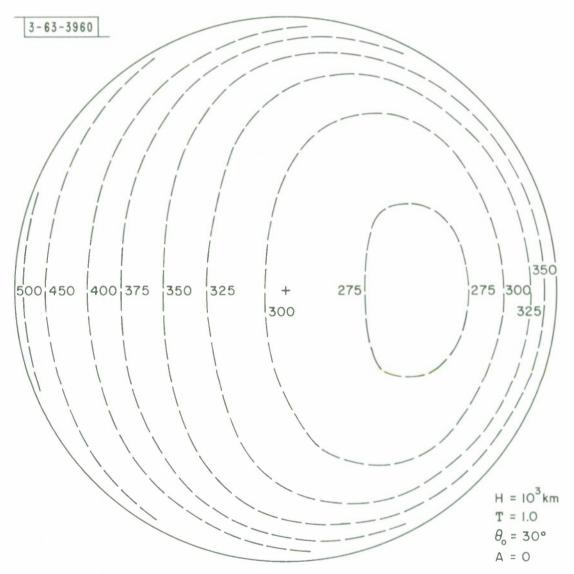


Fig. 46 Map of isophotes for  $H = 10^3$  km, A = 0

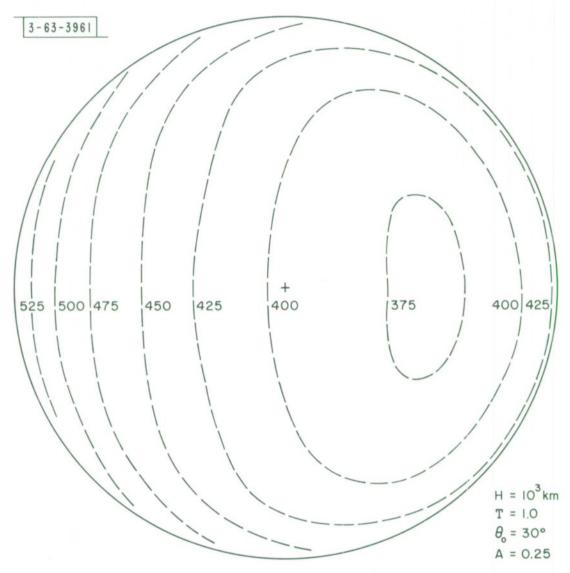


Fig. 47 Map of isophotes for  $H = 10^3$  km, A = 0.25

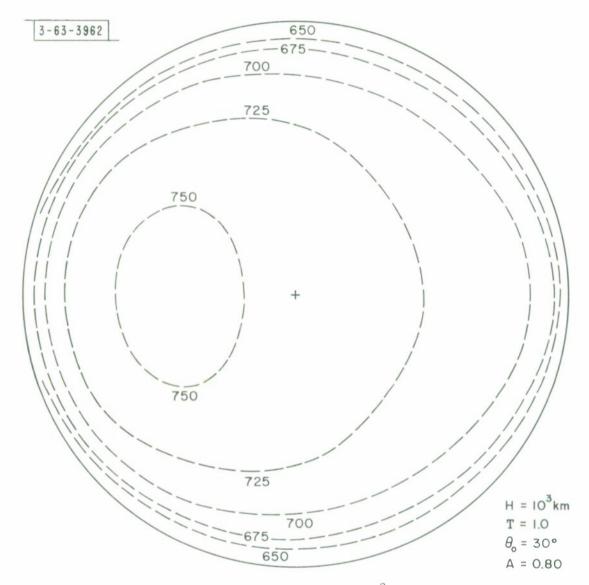


Fig. 48 Map of isophotes for  $H = 10^3$  km, A = 0.80

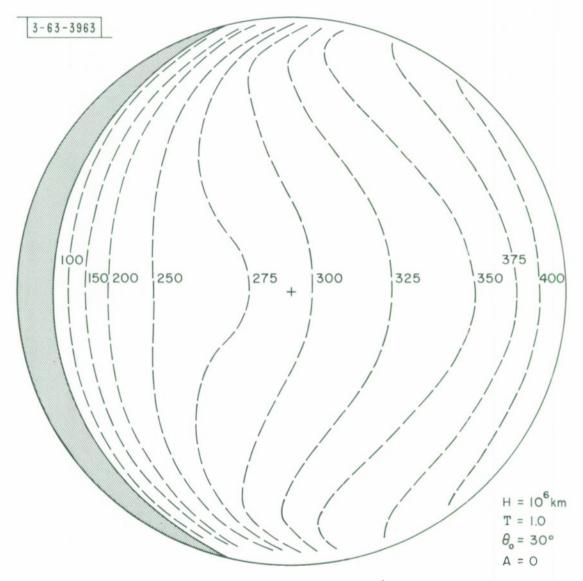


Fig. 49 Map of isophotes for  $H = 10^6$  km,  $_T = 1.00$ 

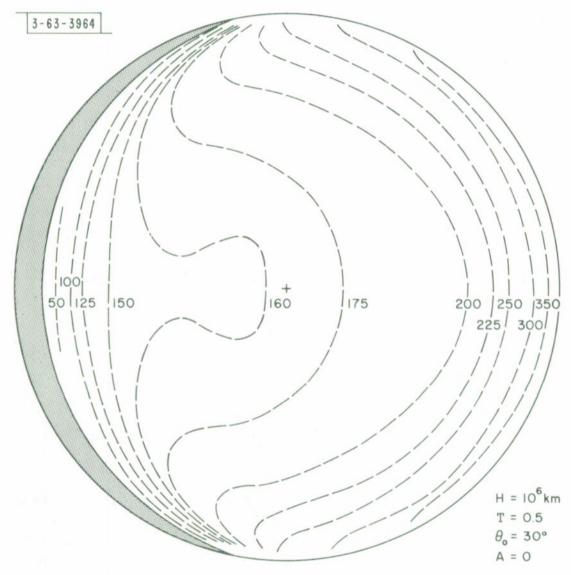


Fig. 50 Map of isophotes for  $H = 10^6$  km,  $\tau = 0.50$ 

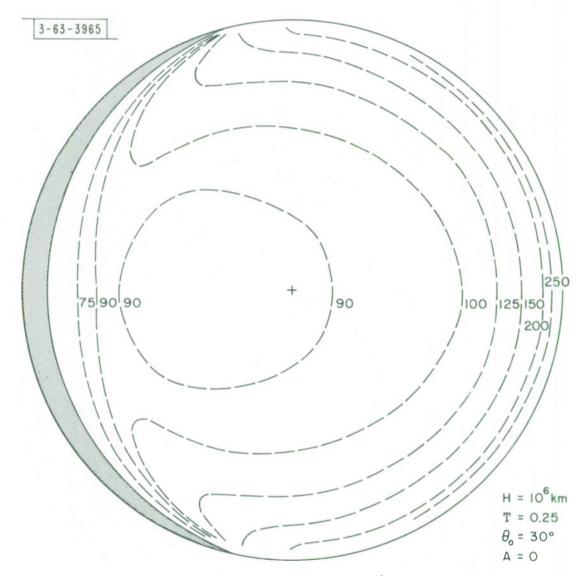


Fig. 51 Map of isophotes for  $H = 10^6$  km,  $\tau = 0.25$ 

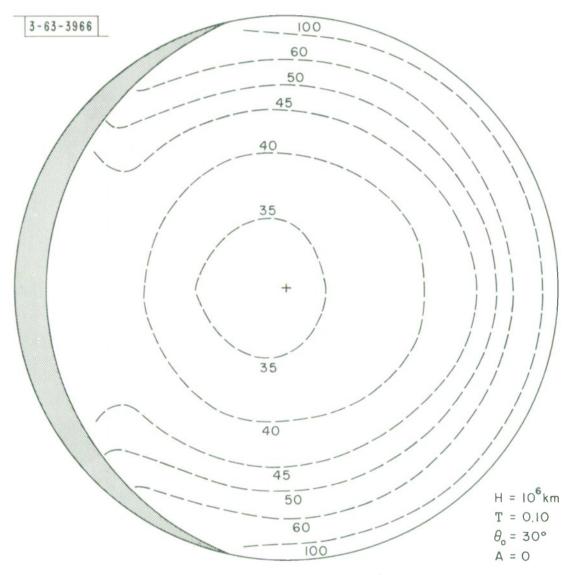


Fig. 52 Map of isophotes for H =  $10^6$  km,  $\tau$  = 0.10

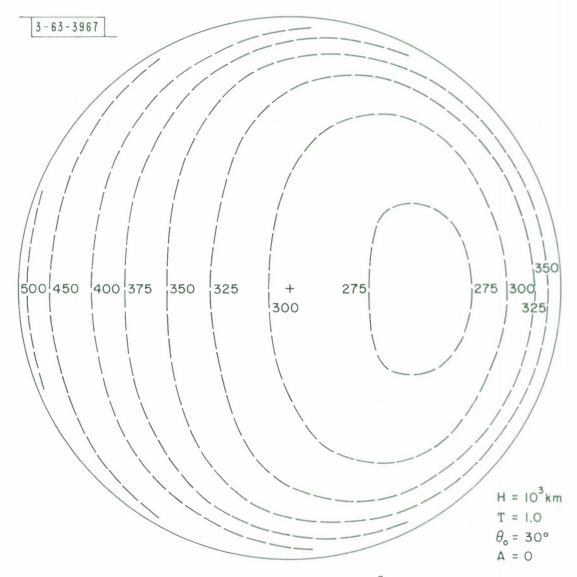


Fig. 53 Map of isophotes for H =  $10^3$  km,  $\tau$  = 1.00

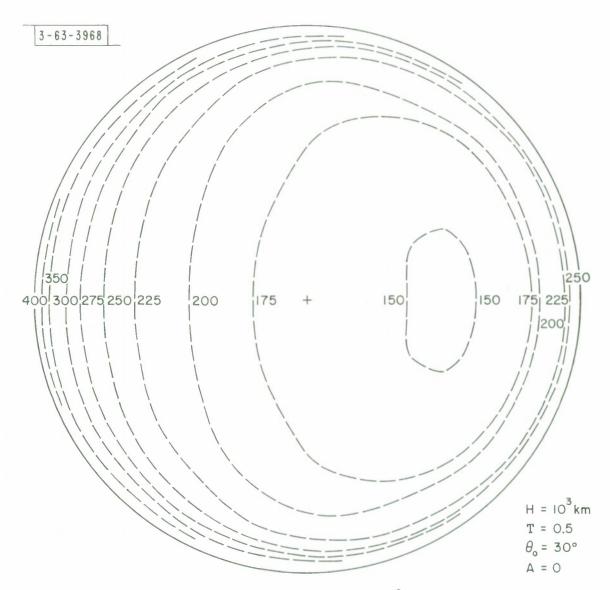


Fig. 54 Map of isophotes for  $H = 10^3$  km,  $\tau = 0.50$ 

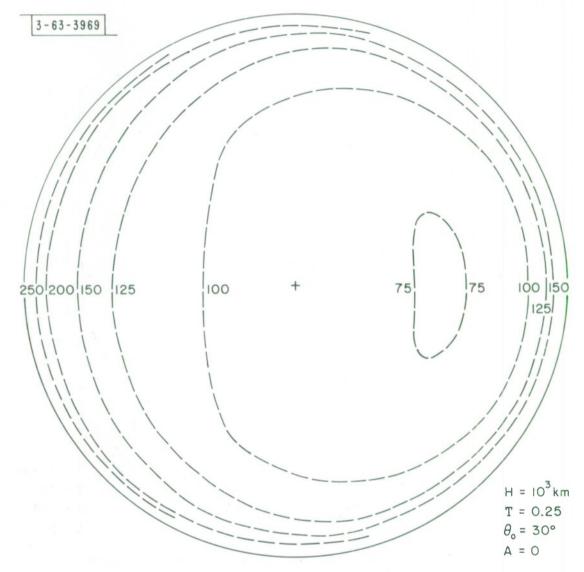


Fig. 55 Map of isophotes for  $H = 10^3$  km,  $\tau = 0.25$ 

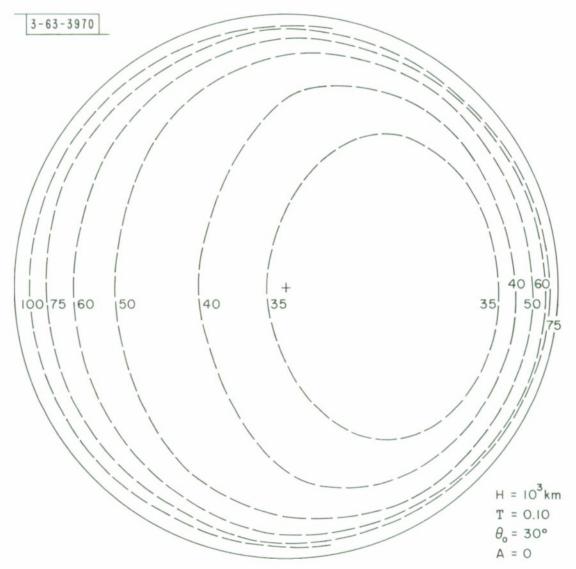


Fig. 56 Map of isophotes for  $H = 10^3$  km,  $\tau = 0.10$ 

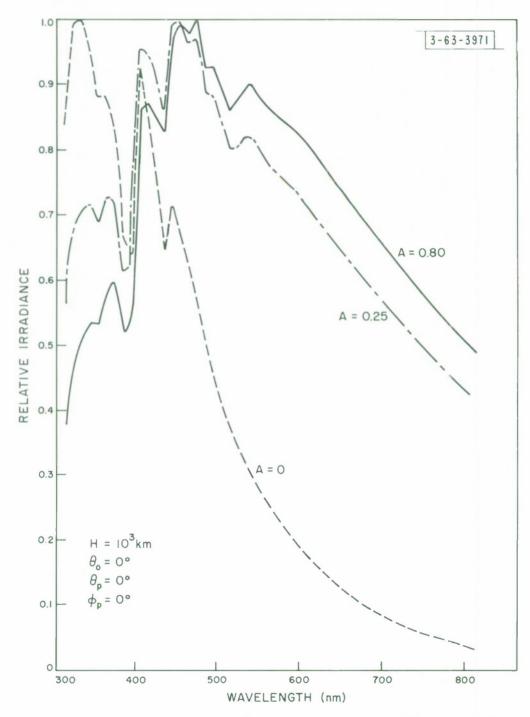


Fig. 57 Relative irradiance for A = 0, 0.25, and 0.80

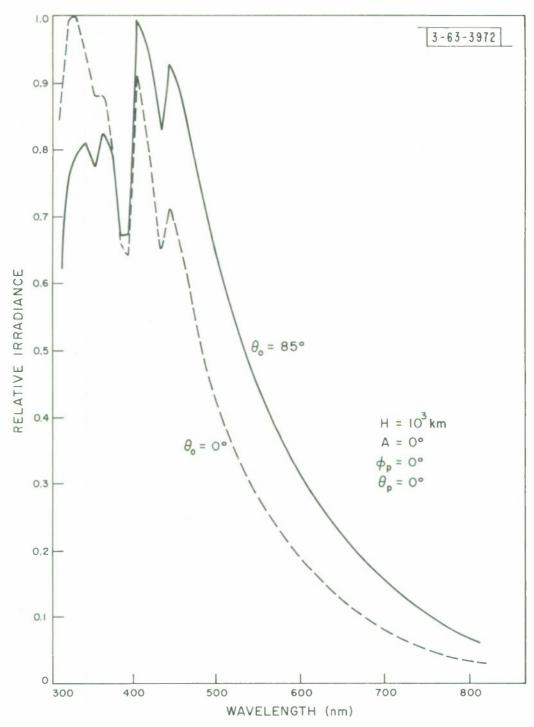


Fig. 58 Relative irradiance for  $\theta_0 = 0$  and  $85^{\circ}$ 

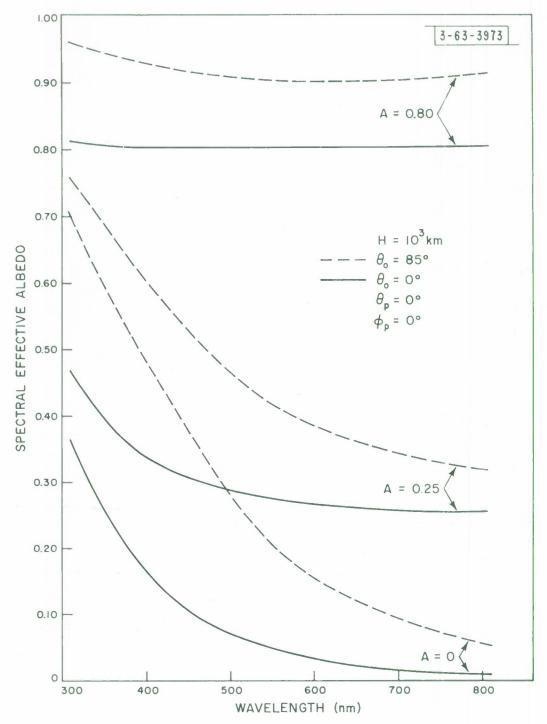


Fig. 59 Spectral effective albedo for A = 0, 0.25 and 0.80;  $\theta_{o} = 0^{o}$  and 85°.

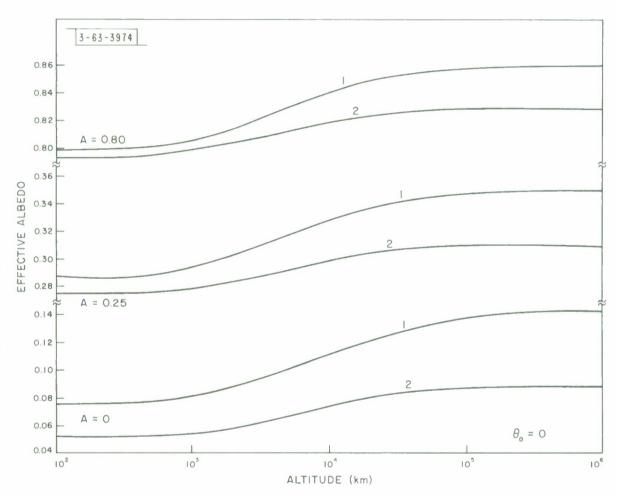


Fig. 60 Effective albedo as a function of altitude for  $\theta_0 = 0^{\circ}$ 

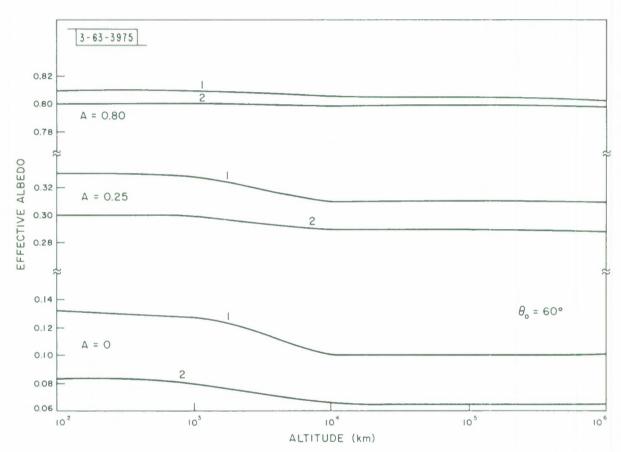


Fig. 61 Effective albedo as a function of altitude for  $\theta_0 = 60^{\circ}$ 

Continuing the type of calculations, the results of which were presented above, the intensity distribution seen by a sensor at synchronous altitude is presented below for specific cases. The numbers have been arrived at by using an extremely useful reference. "Tables Related to Radiation Emerging from a Planetary Atmosphere with Rayleigh Scattering" by Coulson, Dave and Sekera.

In what follows it has been assumed that the sun, the satellite earth sensor, and the center of the earth all lie on a straight line. The symbols used are defined by Figs. 21-22. The sensor sweeps from a spot directly below to the edge of the earth. Corrections have not been made for the curvature of the earth's atmosphere because corrections are only necessary for very oblique angles of incidence (sun zenith angles over 82°) and in such cases the scattered intensity is large (for this geometry).

Figure 62 illustrates the intensity (relative to an incident intensity of 1.0) of scattered light as a function of angle  $\delta$  with wavelength as a parameter. The same information is presented in Fig. 63, but here the intensity is a function of the sensor nadir angle. This makes the sharp increase in intensity near the limb of the earth much more apparent. The increased backscatter near the limb is the reason why a sharp halo of violet rings the earth in the color Gemini IV photographs.

Although the relative intensity of scattered light increases as wavelength decreases, the intensity of sunlight decreases as  $\lambda$  approaches UV wavelengths. The graphs of Fig. 64 have been multiplied by the spectral flux of the sun to obtain the relative intensity of scattered sunlight at various wavelengths.

The data contained in these graphs has been multiplied by the sensor response at two values of  $\delta$  and displayed as a function of  $\lambda$  in Fig. 65. Note that the graph on the right for  $\delta$  = 78.45° is two times larger than the graph for  $\delta$  = 0. The graphs have been put in proper perspective in Fig. 66 where graph (2) is for  $\delta$  = 78.45° graph (3) is for  $\delta$  = 0° and graph (1) is for the intensity distribution over wavelength for a diffusely reflecting atmosphereless earth for  $\delta$  = 0.

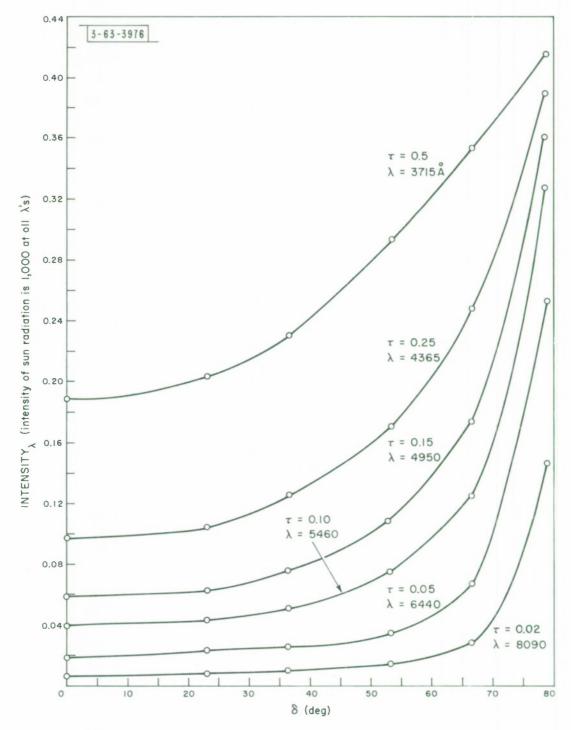


Fig. 62 Relative intensity of scattered light at synchronous altitudes, as a function of the central angle of a great circle ( $\phi = 180^{\circ}$ ,  $\alpha = 180^{\circ}$ ,  $\theta_{o} = 0^{\circ}$ ).

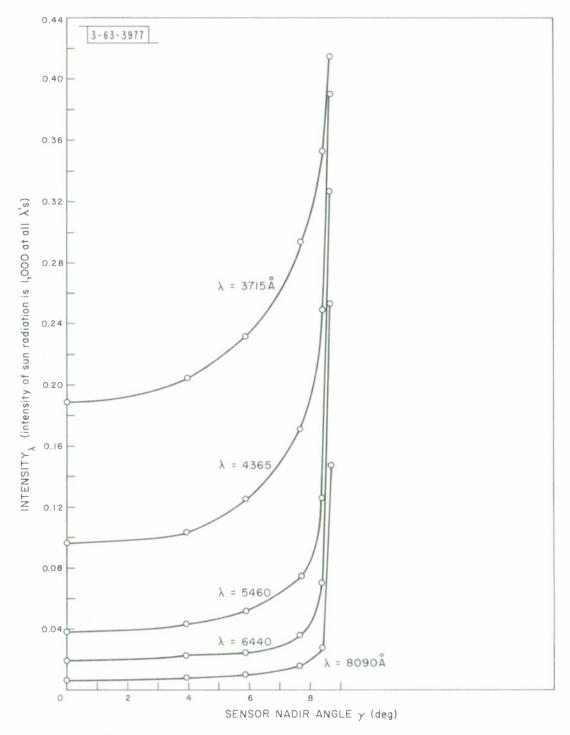


Fig. 63 Relative intensity of scattered light at synchronous altitudes as a function of sensor nadir angle ( $\Phi = 180^{\circ}$ ,  $\alpha = 180^{\circ}$ ,  $\theta = 0^{\circ}$ ).

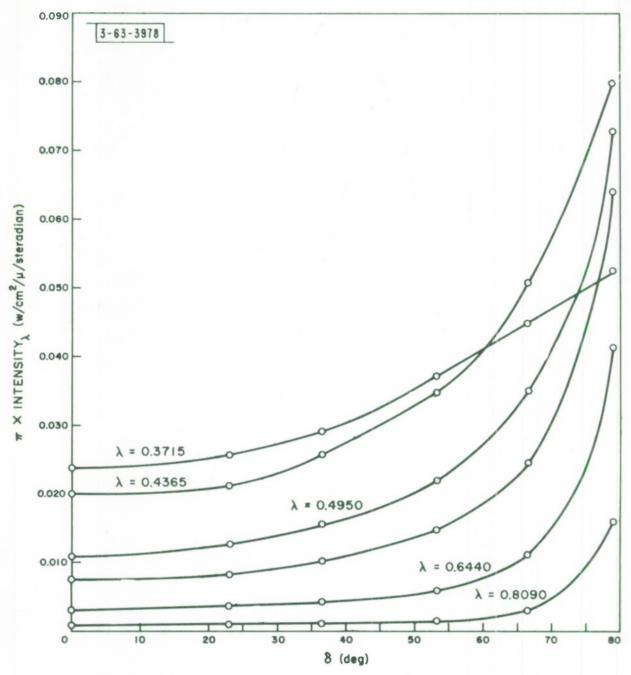


Fig. 64 Relative intensity of scattered sunlight at various wavelengths as a function of the central angle of a great circle ( $\phi = 180^{\circ}$ ,  $\alpha = 180^{\circ}$ ,  $\theta_{o} = 0^{\circ}$ ).

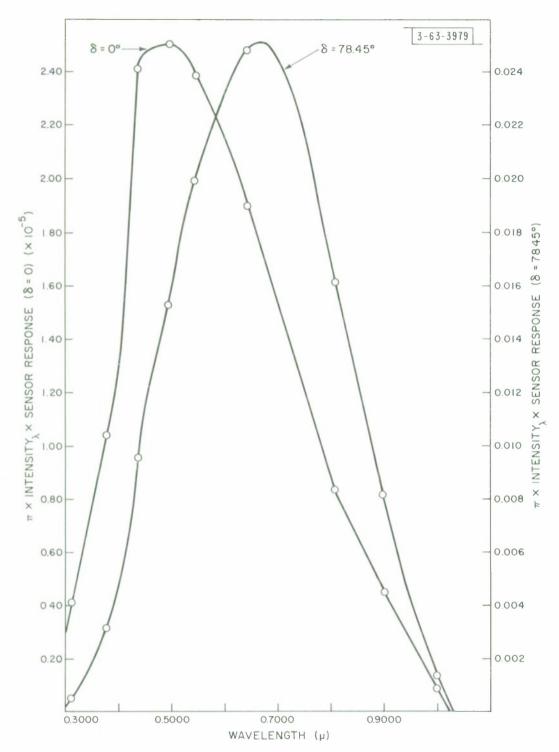


Fig. 65 Light intensity seen by sensor after adjustment for sensor response.

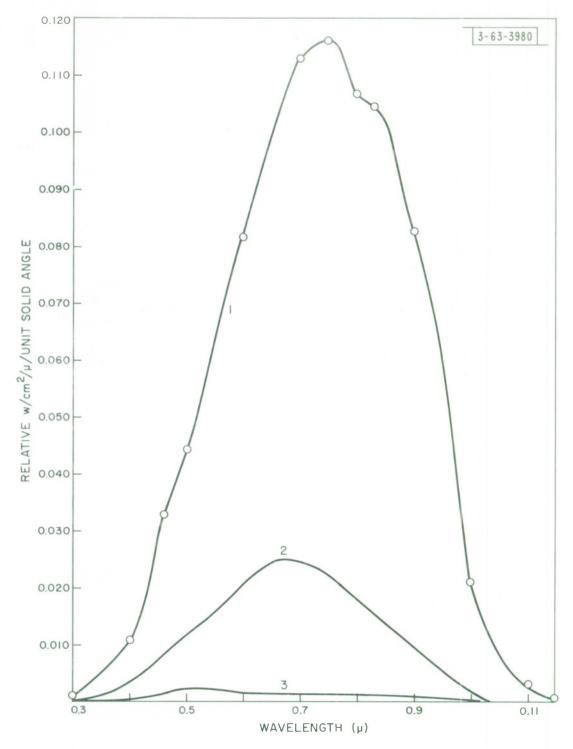


Fig. 66

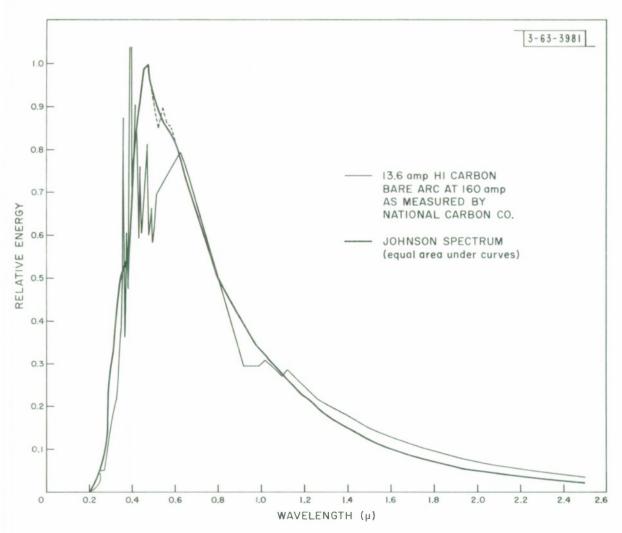


Fig. 67 Relative spectral distribution of the sun

The area under the curves is proportional to the flux seen by the sensor and the current produced. The earth has been assumed flat and the intensity constant in the region viewed. The ratio of the area of graph (2) or (3) to the area of graph (1) is the sensor backscatter coefficient "b<sub>2</sub>".

To predict the current produced by the sensor a very simple experiment can be performed. Measure the current produced by the sensor looking straight down at the light reflected from a diffusely reflecting flat source, illuminated by one solar constant with the same wavelength distribution of intensity as the sun. If this current is  $I_0$  and if the response of the sensor is linear, the current output,  $I_0$ , in outer space for a particular configuration is:

$$I = b_2 I_0$$
  
 $b_2$  for  $\delta = 0.0$  is .0201  
 $b_2$  for  $\delta = 78.45$  is .204

The calculation was done as follows: The geometry of the sun, the earth, and the satellites, is drawn. A value of  $\gamma$  is chosen and the corresponding value of  $\delta$  is gotten from Appendix II. Then  $\gamma + \delta = \cos^{-1}\!\mu$ . The angles  $\cos^{-1}\!\mu_{0}$  and  $\phi$  are gotten from the geometry. The relative intensity can then be found in Table VI of Coulson's tables for all T and hence  $\lambda$ . The value of intensity at a particular  $\lambda$  is multiplied by the flux from the sun at the same wavelength in watts/cm²/micron. (The relative spectral distribution of the sun is given in Fig. 67. The sun's output has been smoothed out so that undue weight is not placed on dips in the curve. This curve must be multiplied by .220 to make the curve absolute.) This value is then multiplied by the relative sensor response. The data points are then plotted versus wavelength in microns and the area under the curve is computed. The ratio of this area to the area under curve (1) in Fig. 66 is the sensor backscatter coefficient "b<sub>2</sub>". This

area is 189.70 boxes or  $474.25 \times 10^{-4}$  sq. units (i. e. each box is  $.05 \times .005 = 25 \times 10^{-4}$  sq. units). Alternately, if the curve (1) is multiplied by the relative sensor response and the area is found the answer is the same but the work is reduced.

Time did not permit additional calculations of  $b_2$  for different satellite orientations. It is felt that the minimum possible value of  $b_2$  for any geometry is not much lower than .02. The calculation, however, does not take into account the effects of water vapor and molecular absorption of radiation at large values of wavelength. Ozone absorption plays a minor role since it becomes important only below 3000 Å. Scattering by dust, haze and aerosols would serve to increase the intensity of scattered light beyond that of a Rayleigh atmosphere.

## APPENDIX I

The Monochromatic Reflectance R $_{\rm V}$  of the Atmosphere-Surface System for Various Values of the Optical Thickness  $_{\rm T}$  and Surface Reflectivity, A.

APPENDIX I

The Monochromatic Reflectance  $R_{\nu}$  of the Atmosphere-Surface System for Various Values of the Optical Thickness  $_{T}$  and Surface Reflectivity, A.

	O Conference Definition A								
τ	λ(A)	μο	0	0.05	0.10	0.15	ectivity, 0.25	0.50	0.80
1.0	3120	0.02 0.10 0.20 0.40 0.60 0.80 0.92 1.00	0.751 0.704 0.650 0.546 0.458 0.392 0.359 0.340	0.758 0.712 0.660 0.558 0.474 0.409 0.377 0.359	0.765 0.721 0.670 0.572 0.490 0.427 0.396 0.378	0.773 0.730 0.681 0.586 0.507 0.446 0.416 0.399	0.790 0.750 0.705 0.616 0.543 0.486 0.469 0.443	0.840 0.809 0.775 0.707 0.651 0.608 0.587 0.575	0.922 0.908 0.891 0.859 0.831 0.811 0.801
0.50	3660	0.02 0.10 0.20 0.40 0.60 0.80 0.92 1.00	0.675 0.610 0.523 0.383 0.297 0.241 0.216 0.202	0.687 0.624 0.540 0.405 0.322 0.268 0.244 0.231	0.698 0.638 0.558 0.428 0.348 0.296 0.273 0.260	0.711 0.653 0.576 0.451 0.374 0.325 0.303 0.291	0.737 0.684 0.614 0.500 0.430 0.385 0.365 0.354	0.809 0.771 0.720 0.638 0.587 0.554 0.540 0.532	0.915 0.898 0.875 0.838 0.816 0.801 0.795 0.791
0.25	4365	0.02 0.10 0.20 0.40 0.60 0.80 0.92 1.00	0.609 0.503 0.375 0.238 0.173 0.136 0.120 0.112	0.625 0.524 0.401 0.270 0.207 0.172 0.157 0.148	0.641 0.545 0.427 0.302 0.242 0.208 0.194 0.186	0.658 0.566 0.454 0.335 0.278 0.245 0.231 0.224	0.693 0.610 0.509 0.402 0.351 0.321 0.309 0.302	0.785 0.727 0.657 0.582 0.546 0.525 0.517 0.512	0.909 0.884 0.854 0.822 0.807 0.798 0.794 0.792
0.15	4950	0.02 0.10 0.20 0.40 0.60 0.80 0.92 1.00	0.569 0.406 0.270 0.158 0.111 0.086 0.076 0.070	0.588 0.432 0.302 0.195 0.151 0.126 0.116 0.111	0.607 0.459 0.335 0.233 0.191 0.167 0.158 0.153	0.627 0.486 0.368 0.271 0.231 0.209 0.200 0.195	0.667 0.541 0.436 0.349 0.313 0.293 0.285 0.281	0.771 0.684 0.612 0.552 0.528 0.514 0.508 0.505	0.905 0.869 0.839 0.814 0.803 0.798 0.796 0.794

Note: The solar zenith angle  $\theta_0$  is given in terms  $\mu_0 = \cos \theta_0$ .

	0		Surface Reflectivity, A							
Τ	λ(A)	μο	0	0.05	0.10	0.15	0.25	0.50	0.80	
0.10	5410	0.20 0.10 0.20 0.40 0.60 0.80 0.92 1.00	0.541 0.323 0.199 0.111 0.077 0.059 0.052 0.048	0.562 0.355 0.236 0.152 0.119 0.102 0.095 0.091	0.583 0.386 0.273 0.193 0.162 0.146 0.139 0.136	0.605 0.418 0.310 0.235 0.205 0.190 0.183 0.180	0.648 0.482 0.386 0.319 0.293 0.279 0.273 0.270	0.760 0.647 0.582 0.536 0.518 0.509 0.505 0.503	0.902 0.855 0.828 0.809 0.802 0.798 0.797 0.796	
0.05	6440	0.02 0.10 0.20 0.40 0.60 0.80 0.92 1.00	0. 475 0. 198 0. 111 0. 059 0. 040 0. 030 0. 026 0. 024	0.352 0.136 0.094 0.072 0.065 0.061 0.059 0.059	0.385 0.180 0.141 0.120 0.113 0.110 0.108 0.107	0. 419 0. 225 0. 188 0. 168 0. 162 0. 158 0. 157 0. 156	0.486 0.315 0.282 0.265 0.259 0.256 0.255 0.254	0.656 0.541 0.519 0.508 0.504 0.501 0.501	0.862 0.815 0.807 0.802 0.800 0.799 0.799	
0.02	8090	0.02 0.10 0.20 0.40 0.60 0.80 0.92 1.00	0.319 0.091 0.048 0.024 0.016 0.012 0.011	0.352 0.136 0.094 0.072 0.065 0.061 0.059 0.059	0.385 0.180 0.141 0.120 0.113 0.110 0.108 0.107	0. 419 0. 225 0. 188 0. 168 0. 162 0. 158 0. 157 0. 156	0.486 0.315 0.282 0.265 0.259 0.256 0.255 0.254	0.656 0.541 0.519 0.508 0.504 0.501 0.501	0.862 0.815 0.807 0.802 0.800 0.799 0.799	
0.01	9600	0.02 0.10 0.20 0.40 0.60 0.80 0.92 1.00	0.197 0.048 0.024 0.012 0.008 0.006 0.005 0.005	0.237 0.095 0.073 0.061 0.057 0.055 0.055	0.277 0.142 0.121 0.110 0.107 0.105 0.104 0.104	0.317 0.189 0.170 0.159 0.156 0.154 0.153 0.153	0.396 0.284 0.267 0.257 0.254 0.253 0.252 0.252	0.597 0.521 0.510 0.504 0.502 0.501 0.500 0.500	0.838 0.808 0.803 0.801 0.800 0.800 0.800 0.799	
0.003	12960	0.02 0.10 0.20 0.40 0.60 0.80 0.92 1.00	0.068 0.013 0.006 0.002 0.001 0.001 0.000 0.000	0.115 0.063 0.056 0.052 0.051 0.050 0.050 0.050	0.161 0.112 0.105 0.102 0.101 0.100 0.100 0.100	0.208 0.161 0.155 0.152 0.151 0.150 0.150 0.150	0.301 0.260 0.254 0.251 0.250 0.250 0.250 0.250	0.534 0.506 0.502 0.500 0.500 0.500 0.500 0.500	0.813 0.802 0.801 0.800 0.800 0.800 0.800 0.800	

	ο λ(A)		Surface Reflectivity, A							
Т		$\mu_{o}$	0	0.05	0.10	0.15	0.25	0.50	0.80	
0.001	17000	0. 02 0. 10 0. 20 0. 40 0. 60 0. 80 0. 92 1. 00	0.023 0.004 0.002 0.000 0.000 0.000 0.000	0.072 0.054 0.051 0.050 0.050 0.050 0.050 0.050	0.121 0.104 0.101 0.100 0.100 0.100 0.100 0.100	0.170 0.153 0.151 0.150 0.150 0.150 0.150 0.150	0. 267 0. 253 0. 251 0. 250 0. 250 0. 250 0. 250 0. 250	0.512 0.502 0.501 0.500 0.500 0.500 0.500 0.500	0.805 0.801 0.800 0.800 0.800 0.800 0.800 0.800	

### APPENDIX II

Angles  $\gamma$  and  $\delta$  Defined by Fig. 21 for the Radius of the Earth = 6371 km and the Radius of the Satellite Orbit = 42,160 km. Angles are Given in Degrees.

Ø.8431864E Ø1 Ø.8080846E Ø2 Ø.8421372E 31 0.8055846E 02 0.8411502E 01 0.8030846E 02 2.8692198E 21 0.6680846E 02 Ø.8401446E 01 Ø.8005846E 02 Ø.8689445E 21 Ø.6655846E Ø2 Ø.8391211E 01 Ø.663Ø846E Ø2 0.7980846E 02 Ø.8688523E 21 0.8380794E Ø1 Ø.7955846E Ø2 Ø.8687432E 21 Ø.66Ø5846E Ø2 0.8370197E 21 Ø.8686172E @1 0.6580846E 02 Ø.793Ø846E Ø2 Ø.8359417E 21 Ø.6555846E Ø2 Ø.79Ø5846E 02 Ø.8684743E 21 8.8348456E 01 Ø.653Ø846E Ø2 Ø.788Ø846E @2 @.8683144E 01 Ø.8337313E 21 Ø.6505846E Ø.7855846E 02 2.8681375E 02 8.8325987E Ø.783Ø846E 02 €.8679436E 0.6488846E 02 01 Ø.8314482E 01 0.8302790E Ø.7805846E 02 Ø.8677326E 21 Ø.6455846E 92 21 Ø.778Ø846E Ø2 0.8675845E 21 Ø.643Ø846E Ø2 Ø.8290918E Ø1 Ø.8672593E ₽1 Ø.7755846E Ø2 0.6405846E 02 €.8278863E Ø1 Ø.773Ø846E Ø2 2.8669972E 91 Ø.638Ø846E Ø2 Ø.8266626E Ø1 0.7705846E 02 2.8667175E 21 Ø.6355846E Ø2 0.8254286E 01 Ø.7680846E Ø2 Ø.8664208E 01 Ø.633Ø846E Ø2 Ø.8241603E Ø.7655846E 0.2 Ø. 8661868E 81 Ø.6305846E Ø2 0.8228817E 0.8657756E 01 Ø.628Ø846E Ø2 0.7630846E 82 Ø.8215849E 21 Ø.76Ø5846E 82 Ø.8654271E @1 Ø.6255846E Ø2 Ø.8292697E 01 Ø.758Ø846E 0.8652613E 01 Ø.623Ø846E Ø2 02 Ø.8189362E 01 €.8646781E Ø.7555846E 0.6205846E 02 01 Ø.8175844E 02 91 Ø.753Ø846E Ø.8642775E 0.6180846E 02 01 92 ₽.8162143E @ 1 Ø.75Ø5846E @.8638595E Ø.6155846E 92 02 01 Ø.8148258E 0.7480846E 02 Ø.8634242E 61 Ø.6130846E 02 C.8134192E Ø.7455846E ₽.8629713E Ø.6105846E 22 02 91 Ø.8119935E 31 3.7430846E 02 9.8625212E 21 Ø.6080846E 22 0.8105504E 31 Ø.6055846E 32 0.7485846E 02 2.8620132E 21 2.8292886E 91 9.7386846E 02 0.6030846E 02 Ø.8Ø76Ø84E Ø1 @.8615277E 21 0.6005846E 02 Ø.7355846E Ø2 €.8609848E 21 0.8061099E 01 Ø.598Ø846E Ø2 Ø.7330846E 22 Ø.86Ø4443E 31 0.8045932E C1 0.5955846E 02 0.8598862E 01 Ø.73Ø5846E Ø2 0.8032578E 21 Ø.593Ø846E Ø2 Ø.728Ø846E Ø2 Ø.8593124E 21 0.8015042E 01 Ø.5905846E Ø2 Ø.7255846E Ø2 2.8587172E 21 0.7999322E 01 0.5880846E 02 Ø.723Ø846E Ø2 2.8581259E 21 Ø.7983419E @1 Ø.7205846E Ø2 0.8574772E 21 Ø.5855846E Ø2 Ø.7967332E 21 Ø.583Ø846E Ø2 Ø.718Ø846E Ø2 Ø.85683Ø5E 31 Ø.7951261E 01 Ø.7155846E 0.5805846E 02 0.7934607E 02 Ø.8561662E 01 21 Ø.578Ø846E Ø2 Ø.713Ø846E 0.7917970E 22 Ø.8554842E 21 01 Ø.7195846E 0.8547843E 3.5755846E 22 21 Ø.7901149E 22 Ø.573Ø846E Ø2 0.7080846E 02 0.8540666E 01 Ø.7884144E Ø.57Ø5846E Ø2 0.7055846E Ø.8533311E 01 2.7866956E 82 0.1 0.7030846E 0.5680846E 02 32 Ø.8525777E 21 Ø.7849585E 91 Ø.7005846E @.8518265E 01 Ø.5655846E 02 €.7832€3£E 02 01 0.6980846E 02 9.8512173E 21 Ø.563Ø846E 912 Ø.7814292E 21 Ø.6955846E Ø2 Ø.5605846E 02 0.8502102E 01 Ø.7796372E 01 0.6930846E 02 Ø.8493853E Ø1 Ø.5580846E Ø2 Ø.7778266E 21 Ø.5555846E Ø2 0.6905846E 02 2.8485423E Ø1 Ø.7759978E 01 Ø.553Ø846E Ø2 Ø.688Ø846E 92 Ø.8476814E Ø1 Ø.77415Ø7E 01 Ø.5505846E 02 Ø.6855846E Ø2 8.8468824E 21 ₽.7722853€ €1

Ø.548Ø846E Ø2

Ø.5455846E Ø2

0.7704817E 01

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0.5430846E	02	Ø.7665795E	Ø 1	Ø.4Ø8Ø846E	Ø2	Ø.6362944E	91
Ø.54Ø5846E	02	Ø.7646411E	21	Ø.4Ø55846E	02	Ø.6334848E	
Ø.538Ø846E	Ø2	Ø.7626843E	91	Ø.4030846E	22	Ø.6324972E	
Ø.5355846E		2.7607094E		Ø.4005846E	Ø2	Ø.6275733E	
Ø.533Ø846E		Ø.7587163E		Ø.398Ø846E	Ø2	Ø.6246332E	
Ø.53Ø5846E	22	2.7567249E	@ 1	Ø.3955846E	22	Ø.6216765E	
Ø.528Ø846E	02	Ø.7546753E	Ø 1	Ø.393Ø846E	02	Ø.6187034E	01
Ø.5255846E	02	2.7526276E	Ø 1	Ø.39Ø5846E	02	0.6157138E	21
Ø.523Ø846E	02	Ø.75Ø5617E	21	0.3880846E	02	Ø.6127282E	
	02	Ø.7484777E	01	Ø.3855846E	02	€.689685SE	
Ø.518Ø846E	02	Ø.7463754E	g1	Ø.383Ø846E	02	Ø.6066475E	
	02	Ø.7442552E	Ø 1	Ø.38Ø5846E	02	0.6035932E	21
	02	Ø.7421168E	Ø 1	Ø.378Ø846E	02	Ø.6005223E	01
		0.7399603E	21	Ø.3755846E	Ø2	Ø.5974356E	
		Ø.7377857E	Ø 1	Ø.373Ø846E	02	Ø.5943329E	21
	02	Ø.7355932E	21	Ø.37Ø5846E	02	@.5912142E	21
	02	Ø.7333826E	91	Ø.368Ø846E	02	Ø.5882796E	01
		Ø.7311540E	91	Ø.3655846E	02	Ø.5849293E	01
		€.7289Ø74E	£ 1	Ø.363Ø846E	32	Ø.5817631E	@ 1
	02	Ø.7266429E	31	Ø.36Ø5846E	02	Ø.5785813E	21
		Ø.7243684E		Ø.358Ø846E	02	Ø.5753837E	01
		0.7220600E		Ø.3555846E	Ø2	0.5721707E	21
	~ ~	9.7197418E		Ø.353Ø846E	22	Ø.5689421E	61
	-	Ø.7174Ø56E		Ø.35Ø5846E	@2	Ø.5656980E	01
	-		01	Ø.348Ø846E	02	Ø.5624386E	21
		Ø.7126799E	Ø1	Ø.3455846E	Ø2	Ø.5591638E	21
	~ ~		21	Ø.343Ø846E	02	Ø.5558738E	21
	~ ~	Ø.7078830E		9.3405846E	02	₽.5525686E	21
	~ ~	0.7054580E		Ø.338Ø846E	02	Ø.5492483E	21
		Ø.7Ø3@153E		Ø.3355846E	Ø2	₽.5459129E	21
			01	Ø.333Ø846E	Ø2	0.5425625E	
Ø.4655846E	~ ~	Ø.6980769E	91	Ø.33Ø5846E	02	0.5391973E	01
	7 -	Ø.6955812E		Ø.328Ø846E	Ø2	Ø.5358172E	
		Ø.693Ø682E		Ø.3255846E Ø.323Ø846E	02	Ø.5324223E	
	~ 0	0.6905372E		0.3205846E	02	Ø.529Ø128E	
			01	Ø.318Ø846E	Ø2	Ø.5255886E	21
	-	Ø • 6854232E	01	Ø.3155846E	Ø2 Ø2	Ø.5221499E	21
		Ø.6828400E Ø.6802394E	91	Ø.3130846E	82	Ø.5186967E	21
	// A	0.6776215E		Ø.31Ø5846E	62	Ø.5152291E	
Ø.443Ø846E		Ø.6749862E	the factoring state of the contract	Ø.3080846E		Ø.5117472E	
	_	Ø.6723336E		Ø.3Ø55846E		@.5Ø82511E	01
		Ø.6696637E		Ø.3Ø3Ø846E		0.5047468E 0.5012164E	
		Ø.6669767E		Ø.3ØØ5846E	02	Ø.497678£E	21
Ø.433Ø846E		0.6642724E		Ø.298Ø846E		2.4941257E	
		Ø.6615511E		_	Ø2	0.4941257E	21
Ø.428Ø846E	0.4	Ø.6588126E		Ø.293Ø846E	02	Ø.4869796E	21
		Ø.656Ø572E			Ø2	2.4833861E	01
		Ø.6532845E		Ø.288Ø846E	02	Ø.4797789E	
		Ø.6504950E		Ø.2855846E		Ø.4761582E	
Ø-4180846E		2.6476885E		Ø.283Ø846E		Ø.4725241E	01
Ø.4155846E		0.6448652E		8.2885846E	02	8.4688766E	-
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Ø.268Ø846E	Ø2	Ø.4541552E Ø	_	Ø.1355846E	Ø2	<b>8.</b> 2377984E Ø1
Ø.2655846E	Ø2	8.4584424E 8			-	Ø.23354Ø3E Ø1
Ø.263Ø846E	02	Ø.4467167E &	_	Ø.13Ø5846E	02	Ø.2292752E Ø1
Ø.2695846E	02	Ø.4429784E Ø	-	Ø.128Ø846E	92	Ø.225@033E 01
Ø.258Ø846E		Ø.4392274E Ø		Ø.1255846E	22	Ø.2207246E Ø1
Ø.2555846E	02	Ø.4354639E 2		Ø.123Ø846E	Ø2	Ø.2164394E Ø1
Ø.253Ø846E	02	Ø.4316879E Ø	_	Ø.12Ø5846E	Ø2	Ø.2121476E 21
Ø.25Ø5846E	Ø2	₿.4278996E @		Ø.118Ø846E	92	Ø.2078495E 81
Ø.248Ø846E	Ø2	Ø.424Ø99RE Ø		Ø.1155846E	Ø2	@.2035451E d1
Ø. 2455846E	22	Ø.42Ø2863E Ø		Ø.113Ø846E	Ø2	Ø.1992346E 81
	22	Ø.4164615E 3		Ø.11Ø5846E	Ø2	Ø.1949182E 21
Ø.243Ø846E	Ø2	Ø.4126248E 2		Ø.1080846E	Ø2	0.1905959E 01
Ø. 2405846E	02	Ø.4Ø87761E 3		Ø.1055846E	Ø2	Ø.1862678E 21
Ø-238Ø846E	Ø2	2.4049157E @	_	Ø.1030846E	Ø2	Ø.1819342E 21
Ø.2355846E	Ø2	Ø.4010437E P	_	Ø.1ØØ5846E	92	Ø.1775951E 21
_	Ø2	Ø.397160£E €	1	Ø.9808463E	@1	Ø.1732586E 81
Ø.23Ø5846E		Ø.3932649E 8	1	Ø.9558464E	Øl	0.16890ESE 01
Ø.228Ø846E		Ø.3893584E Ø	_	3.93Ø8463E	21	Ø.1645461E 81
	02	Ø.3854407E 2	_	0.9058464E	Ø 1	Ø.16Ø1864E Ø1
	02	Ø.3815117E 2		3.8888464E	Øl	Ø.1558218E @1
	02	Ø.3775717E 2	_	Ø.8558463E	Øl	Ø.1514525E €1
	Ø2	Ø.37362Ø7E €	1	Ø.83Ø8464E	<b>Ø</b> 1	Ø.1472786E 21
	Ø2	€.3696589E €	1	J.8058464E	Øl	Ø.1427003E 01
	Ø2	Ø.3656863E @	1	Ø.78Ø8463E	Øl	Ø.1383177E 01
	82	Ø.3617031E Ø	1	Ø.7558464E	Øl	Ø.1339309E Ø1
	02	Ø.3577Ø93E Ø	1	Ø.73Ø8464E	Øl	Ø.12954Ø1E Ø1
	02	₽.3537£5£E €	1	Ø.7Ø58463E	Øl	Ø.1251453E 21
	92	Ø.3496925E Ø	1	Ø.68Ø8464E	Øl	Ø.1207467E 21
	Ø2	Ø. 3456657E 2	1	Ø.6558464E	Ø1	Ø.1163445E ∅1
	82	Ø.3416308E Ø		3.6308463E	Øl	Ø.1119388E Ø1
	02	Ø.3375859E Ø		0.6058464E	Ø1	Ø.1075297E Ø1
	Ø2	Ø.3335311E Ø		0.5808464E	Øl	Ø.1031173E Ø1
	Ø2	2.3294665E Ø		3.5558464E	91	Ø.9872182E 00
	Ø2	Ø.3253922E Ø	1	0.5308464E	Ø1	Ø.9428329E ØØ
	Ø2	Ø.3213083E Ø		Ø.5058464E	01	Ø.8986192E 22
	22	Ø.317215@E 2		Ø.48Ø8464E	Øl	Ø. 8543785E 28
	Ø2	Ø.3131124E 2		Ø.4558464E	Øl	Ø.8101116E 00
Ø.178Ø846E Ø.1755846E	02	6.3696862E		Ø.43Ø8464E	Ø1	Ø.76582Ø2E ØØ
Ø.173Ø846E		Ø.3Ø48795E €		Ø.4Ø58463E		Ø.7215055E 30
0.1705846E		Ø.3Ø07495E 2		Ø.3808463E		Ø.6771689E ØØ
Ø.168Ø846E		Ø.2966126E 2		Ø.3558464E		Ø.6328119E ØØ
		0.2924629E 0		Ø.33Ø8463E		Ø.5884355E 00
Ø.1655846E	02	Ø.2883066E 0		3.3058463E		0.5440413E 80
		Ø.2841417E Ø		Ø-28Ø8464E	Ø1	Ø.49963Ø6E-20
	Ø2	Ø.2799684E Ø		Ø.2558463E	Ø1	2.4552247E-20
	02	Ø.2757869E Ø		Ø.23Ø8463E	Ø1	Ø.4187649E-SØ
	Ø2	Ø.2715971E R		Ø.2858464E	Ø1	Ø.3663127E-₫Ø
	Ø2	Ø.2673992E Ø	_	Ø.18Ø8463E		Ø.3218493E-00
	Ø2	2.2631934E 2		Ø.1558463E	Ø1	0.2773761E-00
Ø.1455846E	Ø2	Ø.2589798E @		Ø.13Ø8464E	Ø1	Ø.2328946E-8Ø
Ø.1430846E		Ø.2547585E Ø		Ø.1058463E		Ø.1884058E-00
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ו יוסרסלמדוים	UZ	Ø.2462931E 8	1	Ø.5584636E		Ø.9941271E-Ø1
				Ø.3084633E		Ø.5491085E-01
				Ø.5846386E	-01	0.1040748E-01

#### REFERENCES

(The References with asterisks are those actually referred to in the paper.)

- \*1. Fritz, Sigmund, "Solar Radiant Energy and its Modification by the Earth and its Atmosphere," Compendium of Meteorology, 1951 pp 13-33.
- \*2. Nordberg, W. and collaborators, "Preliminary Results of Radiation Measurements from the Tiros III Meteorlogical Satellite," NASA Technical Note D-1338, May 1962.
- \*3. Viezee, W. and collaborators, "Variations of Satellite Daytime Radiation Data with Viewing Geometry," Contract AF 19(628)-2777, Scientific Report 3, Air Force Cambridge Research Laboratories, Bedford, Mass., June 1964.
- \*4. Viezee, W. and Davis, P. A., "Analysis of Daytime Radiation Data from Tiros IV," Contract AF 19(628)-2777, Scientific Report 4, Air Force Cambridge Research Laboratories, Bedford, Massachusetts, December 1964.
- \*5. Chandrasekhar, S., Radiative Transfer (Clarendon Press, Oxford 1960)
- \*6. Coulson, K. L., 1959 "Characteristics of the Radiation Emerging from the Top of a Rayleigh Atmosphere, Parts I and II. J. of Planetary and Space Sciences, vol. 1, pp 265-276 and pp 277-284.
- \*7. Coulson, K. L., and Furukawa, P. M., "Distributions of Atmospheric Radiative Heating and Cooling," Contract AF 19(604)-5967, Scientific Report 2, SRI Project 2994, Stanford Research Institute, Menlo Park, California, (Nov. 1960)
- \*8. Snoddy, William C., "Irradiation above the Atmosphere Due to Rayleigh Scattering and Diffuse Terrestrial Reflections," Masters Thesis, University of Alabama, 1964.
- \*9. Coulson, K. L., Dave, J. V., and Sekera, Z. "Table Related to Radiation Emerging from a Planetary Atmosphere with Rayleigh Scattering," 1960, University of California Press.
- 10. Bandeen, W. R., "Tiros II Radiation Data User's Manual," Goddard Space Flight Center, NASA, Greenbelt, Maryland (1962)

- II. Bandeen, W. R., et al., "Infrared and Reflected Solar Radiation Measurements from the Tiros II Meteorological Satellite," J. of Geophysical Research, vol. 66, No. 10, Oct. 1961, pp 3169-3185.
- 12. Fritz, Sigmund, "The Albedo of the Planet Earth and of Clouds," J. Meterol., vol. 6, No. 4, pp 277-282, August 1949.
- 13. Hanel, R. A., and Stroud W. G., The Tiros II Radiation Experiment, "NASA Technical Note D-1152, October 1961.
- 14. Johnson, F. S., "The Solar Constant," J. Meterol., vol. 11, pp 431-439, 1954.
- 15. Nicolet, M. 1951, "Sur la Determination du Flux Energitique du Rayonnement Extraterrestre du Soleil, "Archiv. Meteor., Geophys. u Bioklim, Serie B, vol. 3., p. 209.
- 16. Sherr, P. E. and Wexler, R., Operational Use of Tiros Radiation Measurements, "Contract AF 19(628)-4074, Final Report, Air Force Cambridge Research Laboratories, Bedford, Massachusetts, April 1965.
- 17. Viezee, W. and Davis, P. A., Preliminary Examination of Daytime Radiation Data from Tiros III over Cloudy Regions, "Contract AF 19(628)-2777, Scientific Report 1, Air Force Cambridge Research Laboratories, Bedford, Massachusetts, August 1963.
- 18. Viezee, W. and Davis, P. A., "Studies of Daytime Radiation from Tiros," Contract AF 19(628)-2777, Final Report, Air Force Cambridge Research Laboratories, Bedford, Massachusetts, February 1965.
- 19. Wexler, R. and Sherr, R., "Synoptic Analysis of Tiros III Radiation Measurements," Contract AF 19(628)-429, Final Report, Air Force Cambridge Research Laboratories, Bedford, Massachusetts, Jan. 1964.

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th is ea si sp ph	This paper is a study of the study of the study of the study of various surfaces made to survey typical resurcts covered by a Rayleigh at the study distribution over the earth out and sun. A calculation of the study of the stu	rom the earth. The all, and regions of the eaults of measurements of the mosphere is presented in the properties of the	nedo of rth are f Tiros d in gra orient sponse resente	the earth is defi- discussed. An a data. The albed phical form and ations of observe is made for a sp d to enable interes	ned and attempt do of the the inten- er viewed ecific			
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